

Supporting Information

Shedding light on the photophysical properties of Iridium (III) complexes with *N*-heterocyclic carbene ligands from a theoretical viewpoint

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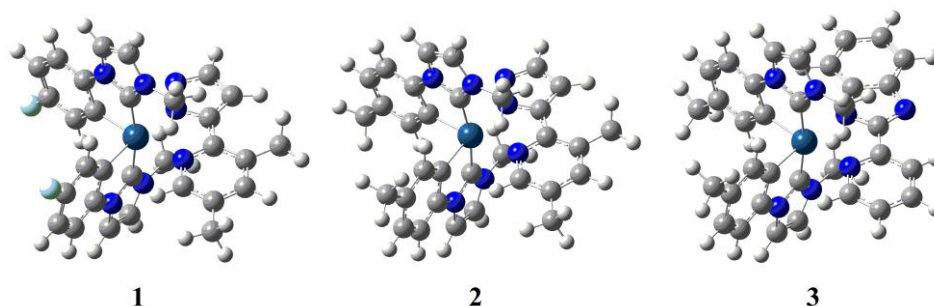


Fig. S1. The optimized structures for complexes 1-3 in the ground state.

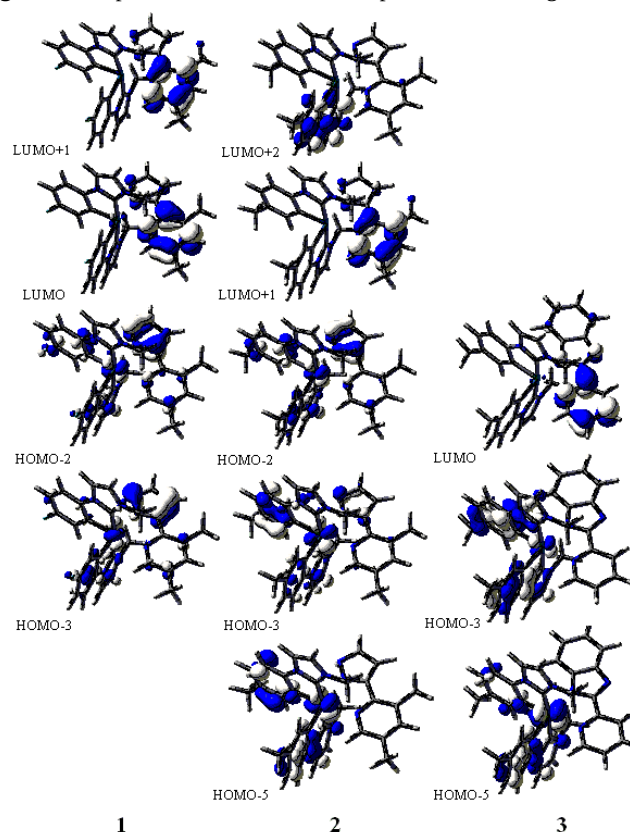


Fig. S2. Diagrams of the related frontier orbitals to the absorption peaks for complexes 1-3 at their optimized S_0 geometry. (Some HOMOs and LUMOs are not shown as they have displayed in the body part of the paper.)

According to Nozaki, when the singlet excited state ($d_{yz} \rightarrow \pi^*$) (denoted as S) couples with the triplet excited state ($d_{xy} \rightarrow \pi^*$) (denoted as T), only the element involving T_y has a nonzero value. The following process will explain how this result comes from.

Usually, the inner product of the angular momentum $\vec{l}$ and the spin of the electron $\vec{s}$ can be represented in the following equation:

$$\vec{l} \cdot \vec{s} = l_x s_x + l_y s_y + l_z s_z = l_z s_z + \frac{1}{2} [l_+ s_- + l_- s_+]$$

Where $l_+ = l_x + il_y$, $l_- = l_x - il_y$. Then, considering the spatial orbital firstly, there will be:¹

$$\begin{aligned} \langle d_{yz} | h_{so} | d_{xy} \rangle &= \langle d_{yz} | l_x + l_y + l_z | d_{xy} \rangle = \langle d_{yz} | l_z + \frac{1}{2} (l_+ + l_-) | d_{xy} \rangle \\ &= \langle d_{yz} | l_z | d_{xy} \rangle + \frac{1}{2} (\langle d_{yz} | l_+ | d_{xy} \rangle + \langle d_{yz} | l_- | d_{xy} \rangle) \\ &= \langle d_{yz}, -2i(d_{x^2-y^2}) \rangle + \langle d_{yz} | l_x | d_{xy} \rangle + \langle d_{yz} | l_y | d_{xy} \rangle \\ &= \langle d_{yz}, -2i(d_{x^2-y^2}) \rangle + \langle d_{yz}, id_{xz} \rangle + \langle d_{yz}, d_{yz} \rangle \end{aligned}$$

Thus, only the matrix element involving the angular momentum of y orientation has a nonzero value. Hence, it is reasonable to take T_y into account in the above mentioned coupling procedure.

The spin functions of the singlet or triplet ones along with the x, y, z coordinate can be written

$$\text{as: } |S\rangle = \frac{1}{\sqrt{2}} |\alpha\beta - \beta\alpha\rangle, \quad |T_x\rangle = \frac{1}{\sqrt{2}} |\beta\beta - \alpha\alpha\rangle, \quad |T_y\rangle = \frac{i}{\sqrt{2}} |\beta\beta + \alpha\alpha\rangle,$$

$$|T_z\rangle = \frac{1}{\sqrt{2}} |\alpha\beta + \beta\alpha\rangle, \text{ respectively.}^2$$

Based on the theory above, considering the spatial orbital and spin coupling together, the complete result will be:³

$$\begin{aligned} \langle T_y | h_{so} | S \rangle &= \zeta \left\langle \frac{i}{\sqrt{2}} C_{d_{yz}} (d_{yz}^\beta + d_{yz}^\alpha) \right| \vec{l} \cdot \vec{s} \left| \frac{1}{\sqrt{2}} C_{d_{xy}} (d_{xy}^\alpha - d_{xy}^\beta) \right\rangle \\ &= \frac{i}{2} \zeta C_{d_{yz}} C_{d_{xy}} \end{aligned}$$

The calculated soc matrix elements are listed in **Table S1**.

Table S1. SOC matrix elements between the $^3(d_k\pi^*)$ and $^1(d_l\pi^*)$ states.^{1, 3-5}

	$^1(d_{xz}\pi^*)$	$^1(d_{yz}\pi^*)$	$^1(d_{xy}\pi^*)$
$^3(d_{xz}\pi^*)$	0	$-\frac{i}{2} C_{xz} C_{yz}$	$\frac{i}{2} C_{xz} C_{xy}$
$^3(d_{yz}\pi^*)$	$\frac{i}{2} C_{yz} C_{xz}$	0	$-\frac{i}{2} C_{yz} C_{xy}$
$^3(d_{xy}\pi^*)$	$-\frac{i}{2} C_{xy} C_{xz}$	$\frac{i}{2} C_{xy} C_{yz}$	0

Table S2. Electronic transitions of **1** calculated at optimized T_1 geometry by TDDFT in CH_2Cl_2 solution. The solvent effect is modeled by the PCM method.^a

S_n/T_m	Energy (cm^{-1})	λ (nm)	Major contribution ^b
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T ₁	18543	539	H→L(0.71), H-1→L(0.49), H-2→L(-0.38)
T ₂	24745	404	H→L(0.37), H-2→L(0.71), H-4→L(-0.38)
S ₁	25410	393	H→L(0.98)
S ₂	27460	364	H-1→L(0.96)
S ₃	29485	339	H-2→L(0.96)
S ₄	31100	321	H-3→L(0.91)
S ₅	33082	302	H→L+1(0.75), H→L+2(0.40), H→L+4 (-0.31)

^a Only the triplet excited states below the first singlet excited states are listed.

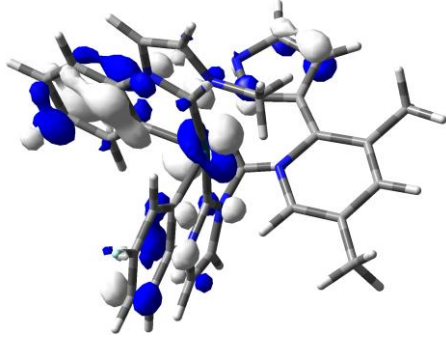
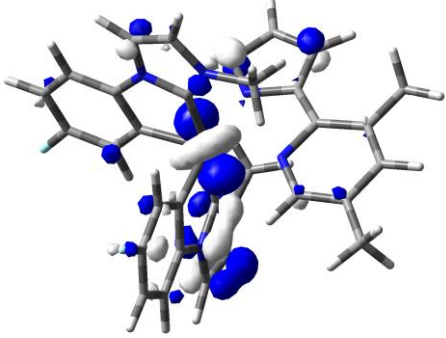
^b the values in the parentheses are the CI coefficients (a_i) of a particular transition. Only the abs of CI coefficients >0.20 are displayed.

Table S3. The Calculated transition dipole moment (unit: a.u.) from the ground state to the singlet excited state together with the oscillator strengths.^a

State	X	Y	Z	f
S ₁	-0.3831	0.6091	0.0705	0.0403
S ₂	-0.8662	0.5987	0.0353	0.0926
S ₃	-0.5351	-1.0390	-0.1475	0.1243
S ₄	0.5266	-0.9062	-0.1891	0.1071
S ₅	0.2207	-0.2672	-0.3967	0.0279

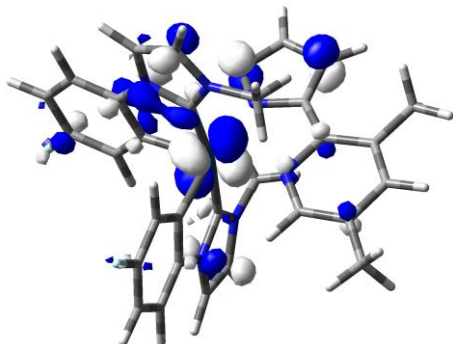
a: Oscillator strengths and transition dipole moment are some intrinsic characteristics of the excited states, here we use the larger one obtained in the TDDFT calculation on respective S₀ and T₁ geometry for the following calculation of the radiative decay rates.

Table S4. MO surface and orbital contributions of **1** at optimized T₁ geometry.

MO	MO surface	Nature(c_d^a)
HOMO-4		d _{xz} (0.3187)
HOMO-3		d _{yz} (-0.3575)

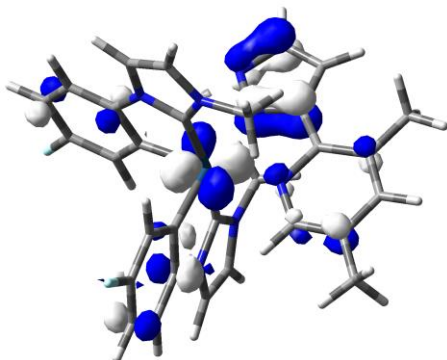
HOMO-2

$d_{xz}(-0.4242)$



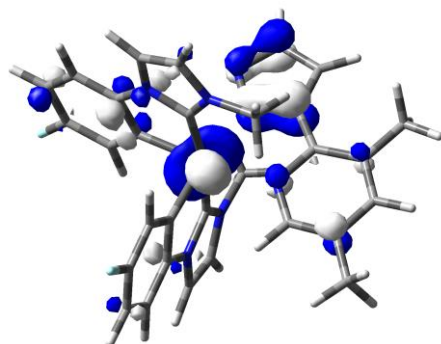
HOMO-1

$d_{xy}(-0.3067)$



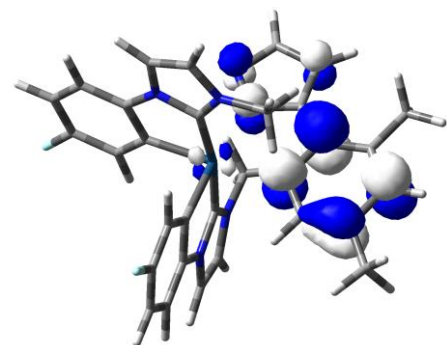
HOMO

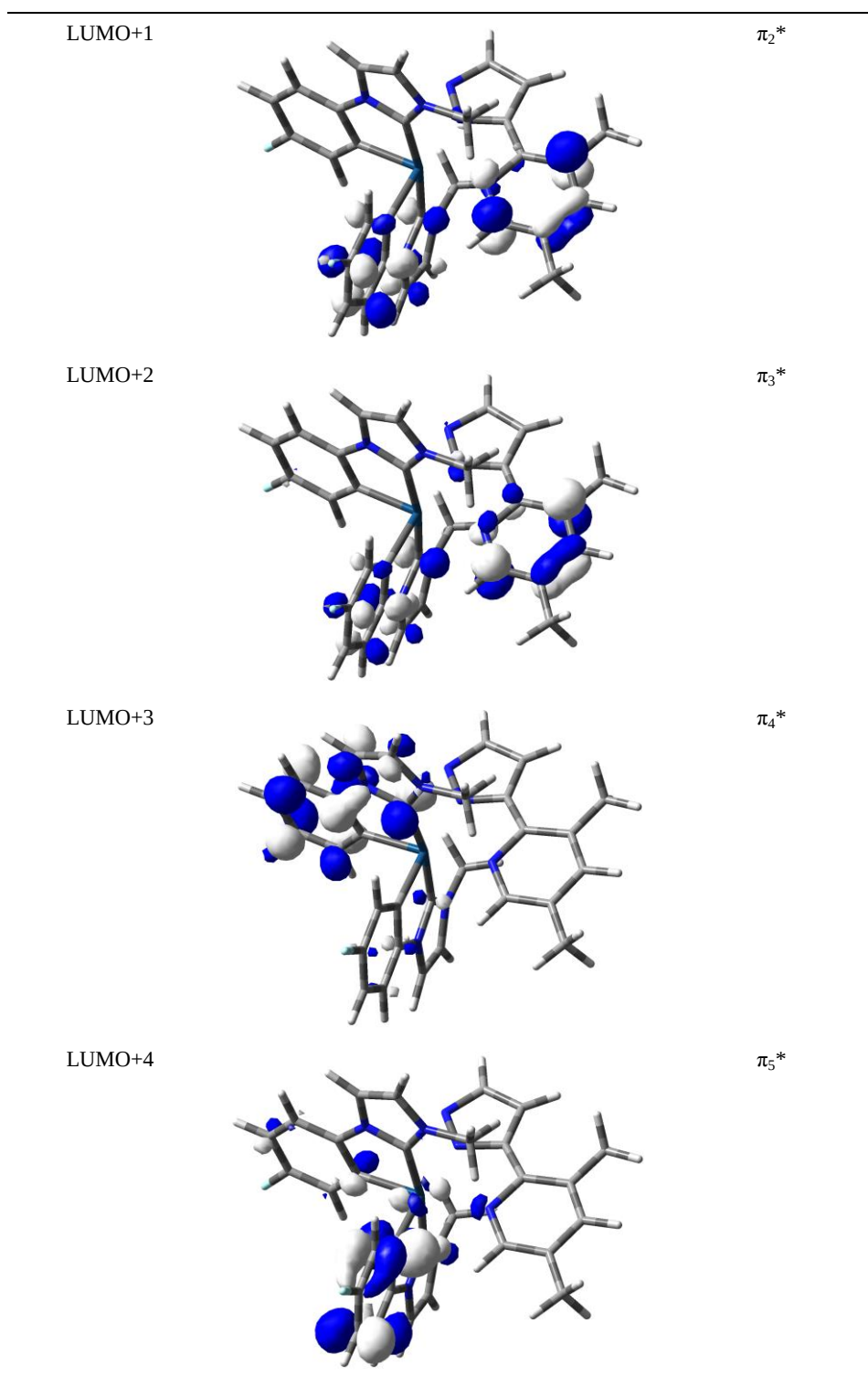
$d_{yz}(-0.3927)$



LUMO

π_1^*





^a Only those with $|c_d| > 0.3$ are given in **Table S4**.

Table S5. SOC matrix elements $\langle T_1^a | H_{\text{soc}} | S_n \rangle$ (cm^{-1}) of **1** calculated at the optimized T_1 geometry. The calculated radiative decay rate constant (k_r/s^{-1}) is also shown in the **Table S5**.

E(T_1)=18543 cm^{-1}									
S_n	$\langle T_1^x H_{\text{soc}} S_n \rangle$	$\langle T_1^y H_{\text{soc}} S_n \rangle$	$\langle T_1^z H_{\text{soc}} S_n \rangle$	$\Delta E(T_1^x)$	$\Delta E(T_1^y)$	$\Delta E(T_1^z)$	$k_r^x/8.478 \times 10^{-30} \text{C} \cdot \text{m}$	$k_r^y/8.478 \times 10^{-30} \text{C} \cdot \text{m}$	$k_r^z/8.478 \times 10^{-30} \text{C} \cdot \text{m}$
S_1		128.85	137.48		-2.418	-2.752		0.0114289	0.0014114
S_2	-105.82	-181.41		-1.256	-3.691		0.010279	-0.01218	
S_3	-137.17		250.91	-1.720		-5.754	0.006708		-0.003382

S ₄	108.80	116.09	-0.943	-1.073	-0.007852	-0.001748
S ₅						
k _r ^a			-2.976	-7.052	-9.579	10767.44
k _r						2761.4107
						516.0841
						4681.6465

The blank represents that there is no nonzero coupling value at the specified coupling states.

Table S6. Electronic transitions of **2** calculated at optimized T₁ geometry by TDDFT in CH₂Cl₂ solution. The solvent effect is modeled by the PCM method.^a

S _n /T _m	Energy (cm ⁻¹)	λ (nm)	Major contributions ^b
T ₁	18559	539	H→L(0.54), H-1→L(0.58), H-2→L(-0.47)
T ₂	24285	412	H→L(0.40), H-1→L(0.40), H-2→L(0.61), H-4→L(0.44)
S ₁	24375	410	H→L(0.99)
S ₂	26666	375	H-1→L(0.96)
S ₃	26900	347	H-2→L(0.95)
S ₄	30593	327	H-3→L(0.89)
S ₅	32200	311	H→L+2(-0.44), H→L+1(0.68), H-4→L(-0.45)

^a Only the triplet excited states below the first singlet excited states are listed.

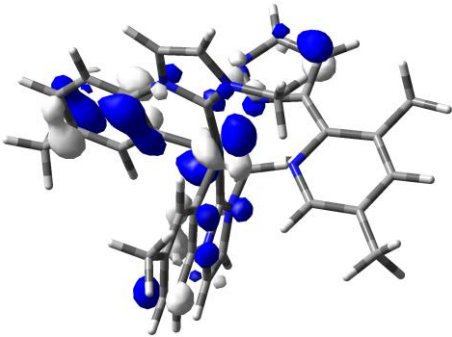
^b the values in the parentheses are the CI coefficients(a_i) of a particular transition. Only the abs of CI coefficients >0.20 are displayed.

Table S7. The Calculated transition dipole moment (unit: a.u.) from the ground state to the excited state together with the oscillator strengths.^a

State	X	Y	Z	f
S ₁	-0.1813	0.4161	0.0196	0.0153
S ₂	-0.8815	0.5989	-0.0214	0.0920
S ₃	-0.6111	-0.8854	-0.0581	0.1017
S ₄	-0.5403	1.0666	0.1091	0.1339
S ₅	-0.0392	0.4607	0.4750	0.0430

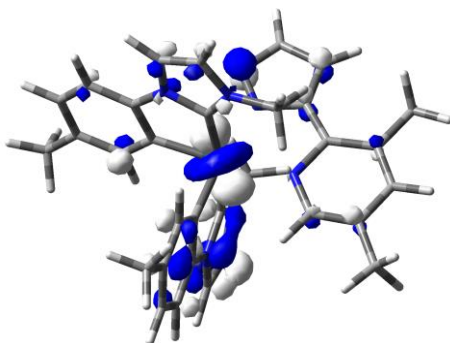
a: Oscillator strengths and transition dipole moment are some intrinsic characteristics of the excited states, here we use the larger one obtained in the TDDFT calculation on respective S₀ and T₁ geometry for the following calculation of the radiative decay rates.

Table S8. MO surface and orbital contributions of **2** at optimized T₁ geometry.

MO	MO surface	Nature(c _d ^a)
HOMO-4		d _{xz} (-0.3542)

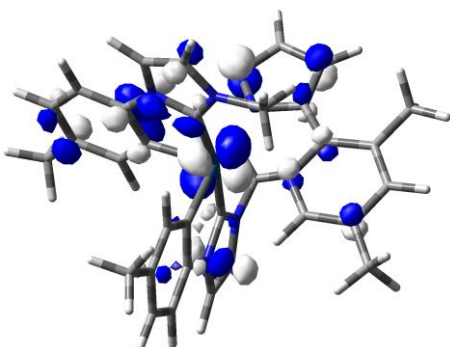
HOMO-3

$d_{yz}(0.3558)$



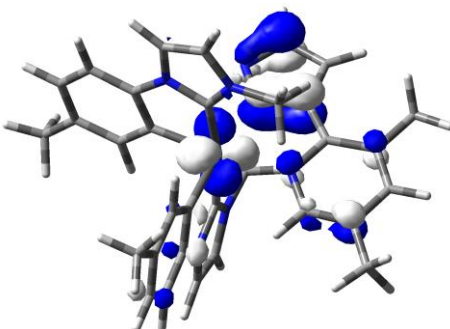
HOMO-2

$d_{xz}(-0.4075)$



HOMO-1

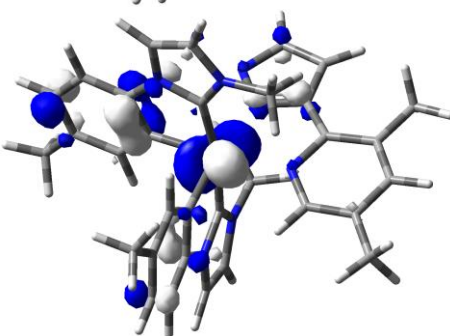
$d_{yz}(-0.3112)$



HOMO

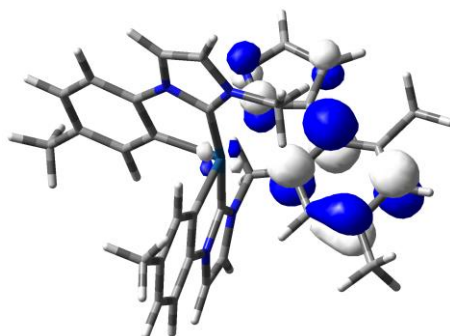
$d_{yz}(-0.3417)$

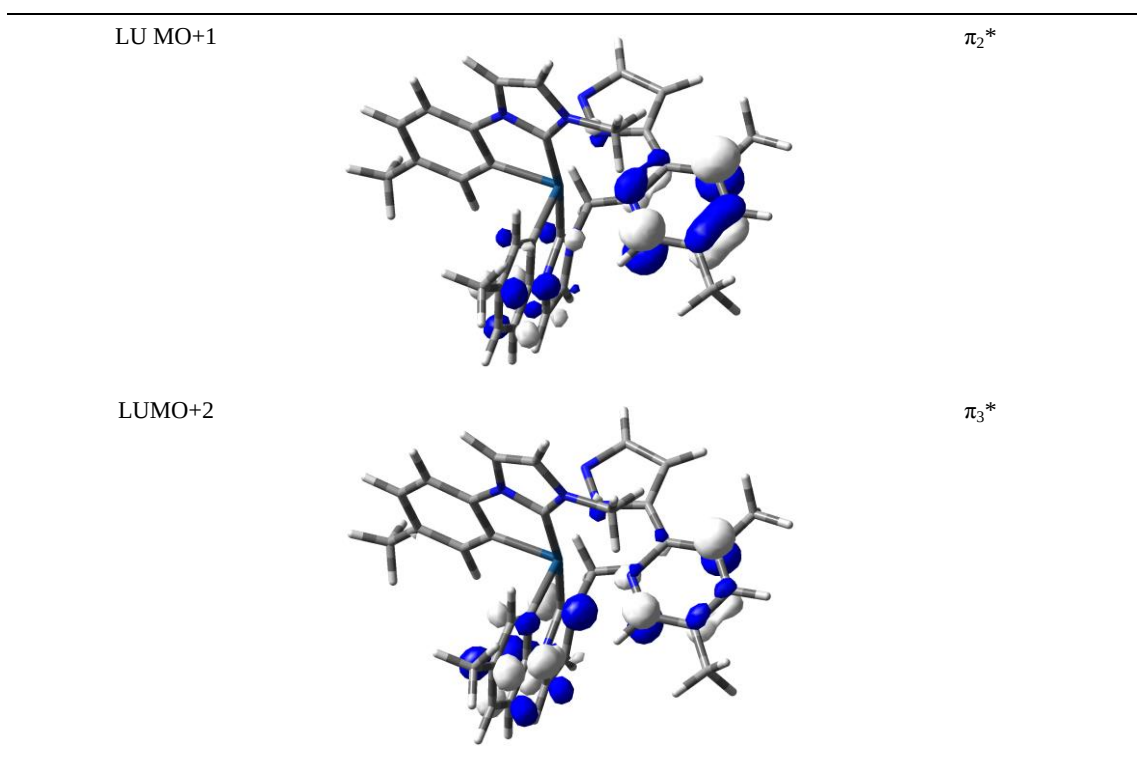
$d_{xy}(0.3733)$



LUMO

π_1^*





^a Only those with $|c_d| > 0.3$ are given in **Table S8**.

Table S9. SOC matrix elements $\langle T_1^a | H_{\text{soc}} | S_n \rangle$ (cm^{-1}) of **2** calculated at the optimized T_1 geometry. The calculated radiative decay rate constant (k_r/s^{-1}) is also shown in the **Table S9**.

E(T_1)=18559 cm^{-1}									
S_n	$\langle T_1^x H_{\text{soc}} S_n \rangle$	$\langle T_1^y H_{\text{soc}} S_n \rangle$	$\langle T_1^z H_{\text{soc}} S_n \rangle$	$\Delta E(T_1^x)$	$\Delta E(T_1^y)$	$\Delta E(T_1^z)$	$k_n^x/8.478 \times 10^{-30} \text{C} \cdot \text{m}$	$k_n^y/8.478 \times 10^{-30} \text{C} \cdot \text{m}$	$k_n^z/8.478 \times 10^{-30} \text{C} \cdot \text{m}$
S_1	155.67	147.70	142.49	-4.167	-3.751	-3.491	-0.00485	0.0105669	0.0004802
S_2		-132.98	126.06		-2.181	-1.960		-0.009823	-0.000333
S_3	171.57		311.37	-3.529		-11.623	-0.01019		-0.001757
S_4		140.86	-133.53		-1.649	-1.482		0.012485	-0.001211
S_5	-71.23		-129.27	-0.372		-1.225	0.000205		-0.004501
k_r^a				-8.068	-7.581	-19.781	8231.553	6547.1485	2005.6726
k_r								5594.7913	

The blank represents that there is no nonzero coupling value at the specified coupling states.

Table S10. Electronic transitions of **3** calculated at optimized T_1 geometry by TDDFT in CH_2Cl_2 solution. The solvent effect is modeled by the PCM method.^a

S_n/T_m	Energy (cm^{-1})	λ (nm)	Major contribution ^b
T_1	16259	615	$\text{H} \rightarrow \text{L}(0.91)$, $\text{H-1} \rightarrow \text{L}(-0.31)$
T_2	20257	494	$\text{H-2} \rightarrow \text{L}(0.45)$, $\text{H-3} \rightarrow \text{L}(0.72)$
S_1	21051	475	$\text{H} \rightarrow \text{L}(0.98)$
S_2	23128	432	$\text{H-1} \rightarrow \text{L}(0.98)$
S_3	26421	378	$\text{H-2} \rightarrow \text{L}(0.92)$, $\text{H-4} \rightarrow \text{L}(-0.28)$
S_4	28463	351	$\text{H} \rightarrow \text{L+1}(-0.44)$, $\text{H-2} \rightarrow \text{L}(0.31)$, $\text{H-3} \rightarrow \text{L}(-0.57)$, $\text{H-4} \rightarrow \text{L}(0.59)$
S_5	28562	350	$\text{H} \rightarrow \text{L+1}(0.83)$, $\text{H-3} \rightarrow \text{L}(-0.52)$

^a Only the triplet excited states below the first singlet excited states are listed.

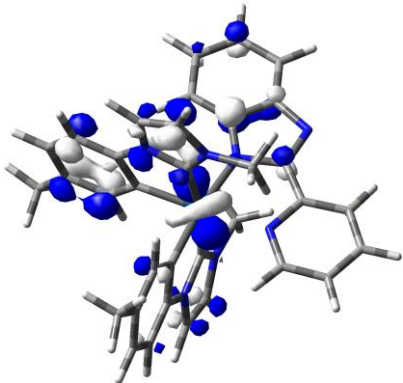
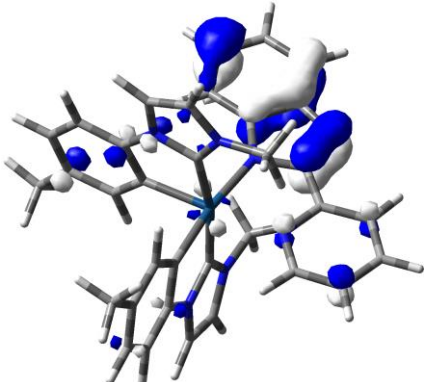
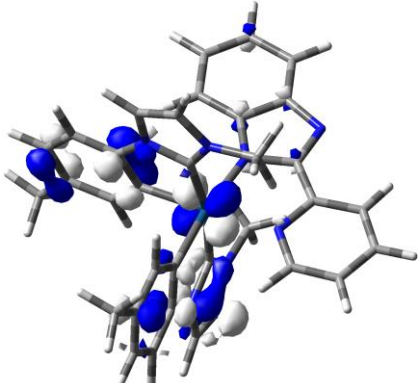
^b the values in the parentheses are the CI coefficients(a_i) of a particular transition. Only the abs of CI coefficients > 0.20 are displayed.

Table S11. The Calculated transition dipole moment (unit: a.u.) from the ground state to the excited state together with the oscillator strengths.^a

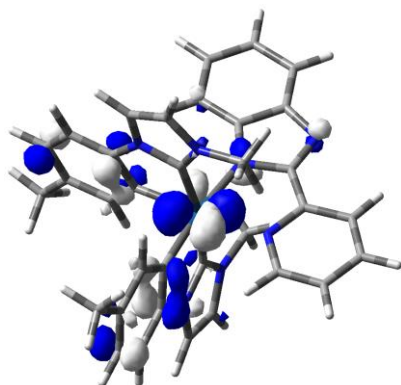
State	X	Y	Z	f
S ₁	0.0986	0.2267	0.0052	0.0046
S ₂	0.2526	-0.8037	-0.0141	0.0573
S ₃	1.5167	0.4464	-0.0003	0.2202
S ₄	-0.0847	0.0802	-0.0836	0.0019
S ₅	-1.3003	-1.1096	0.0552	0.2692

a: Oscillator strengths and transition dipole moment are some intrinsic characteristics of the excited states, here we use the larger one obtained in the TDDFT calculation on respective S₀ and T₁ geometry for the following calculation of the radiative decay rates.

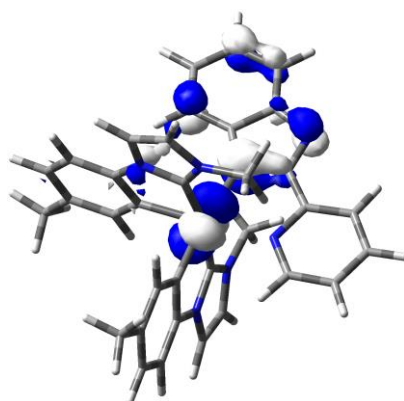
Table S12. MO surface and orbital contributions of **3** at optimized T₁ geometry.

MO	MO surface	Nature(c _d ^a)
HOMO-4		π_1
HOMO-3		π_2
HOMO-2		d _{yz} (0.3793)

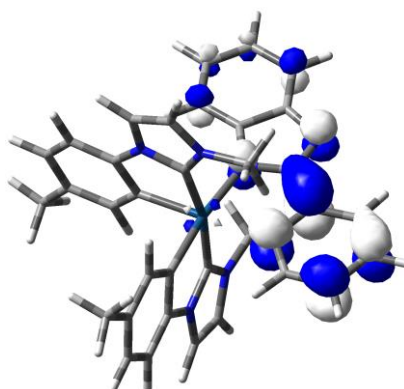
HOMO-1

 $d_{xy}(0.5248)$ 

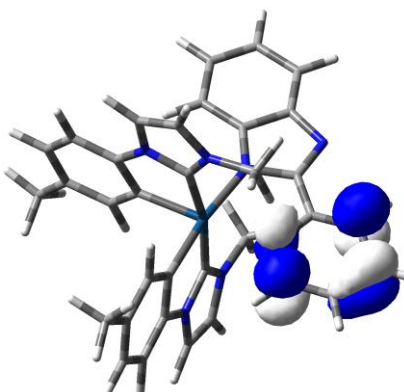
HOMO

 $d_{xz}(0.5091)$ 

LUMO

 π_3^* 

LUMO+1

 π_4^* 

^a Only those with $|c_d| > 0.3$ are given in **Table S12**.

Table S13. SOC matrix elements $\langle T_m^a | H_{\text{soc}} | S_n \rangle$ (cm^{-1}) of **3** calculated at the optimized T_1 geometry. The calculated radiative decay rate constant (k_r/s^{-1}) is also shown in the **Table S13**.

$E(T_1) = 16259 \text{ cm}^{-1}$

S_n	$\langle T_1^x H_{\text{soc}} S_n \rangle$	$\langle T_1^y H_{\text{soc}} S_n \rangle$	$\langle T_1^z H_{\text{soc}} S_n \rangle$	$\Delta E(T_1^x)$	$\Delta E(T_1^y)$	$\Delta E(T_1^z)$	$k_n^x / 8.478 \times 10^{-30} \text{C} \cdot \text{m}$	$k_n^y / 8.478 \times 10^{-30} \text{C} \cdot \text{m}$	$k_n^z / 8.478 \times 10^{-30} \text{C} \cdot \text{m}$
S_1	179.67			-6.737			0.003696		
S_2	522.67			-39.771			0.019219		
S_3		-126.07	-355.86		-1.564	-12.462		-0.005538	0.00001051
S_4		-42.68	-120.46		-0.149	-1.189		-0.00028	0.0008251
S_5									
k_r^a				-46.508	-1.713	-13.651	19595.75	1263.1755	26.056437
k_r								6961.6603	

The blank represents that there is no nonzero coupling value at the specified coupling states.

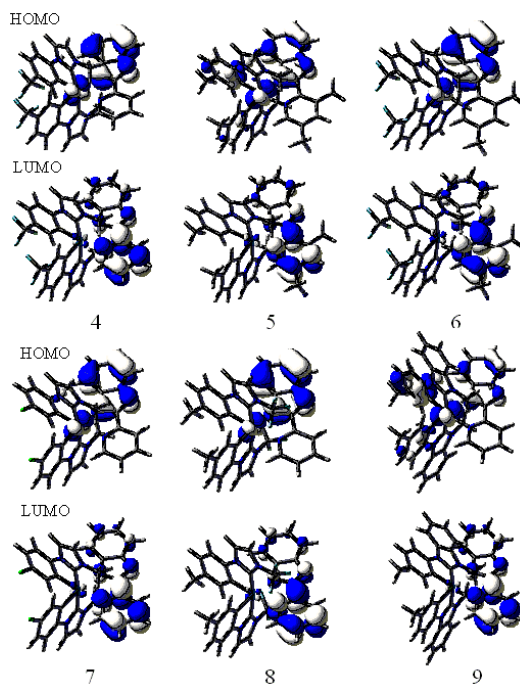


Fig. S3. Diagrams of the related frontier orbitals for complexes **4-9** with different substituents at their optimized S_0 geometry respectively.

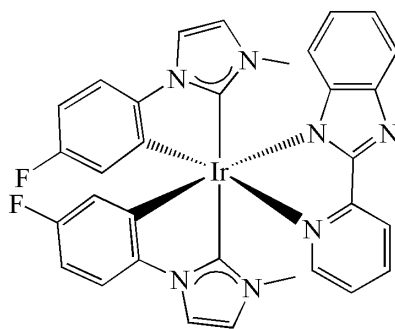


Fig. S4. Molecular graph for complex with F as the substituent at R_1 sites.

Table S14. Calculated frontier molecular orbital energy levels and the energy gap for complex with F as the substituent at R_1 site based on optimized S_0 geometry.

HOMO	LUMO	Energy gap
-5.52	-1.64	3.89

Table S15. NBO analysis for N-C bond between carbene ring and the substituents on it of complexes **3** and **8**.

Complexes	Occupancy	Bond orbital type	Contributions
3	1.9894	BD(1)	64.26%N

8	1.9878	BD(1)	35.74%C
			63.55%N
	1.6322	BD(2)	36.45%C
			97.48%N
			2.52%C

Reference

- [1] C. J. Ballhausen, *Introduction to Ligand Field Theory*, McGraw-Hill, New York, 1969.
- [2] K. Nozaki, *Journal of the Chinese Chemical Society* **2006**, 53, 101.
- [3] S. Obara, M. Itabashi, F. Okuda, S. Tamaki, Y. Tanabe, Y. Ishii, K. Nozaki, M. A. Haga, *Inorg. Chem.* **2006**, 45, 8907.
- [4] G. S. M. Tong, C. M. Che, *Chem. -Eur. J.* **2009**, 15, 7225.
- [5] A. F. Rausch, H. H. H. Homeier, H. Yersin, *Organometallic Pt(II) and Ir(III) triplet emitters for OLED applications and the role of spin-orbit coupling: A study based on high-resolution optical spectroscopy*, Springer Berlin Heidelberg, 2010.