

RANS models deficiencies, palliatives, and corrections

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- The Boussinesq approximation (or EVM) is a brutal simplification of reality, and this can be a major source of predictive defects.
- And this is regardless of the closure approximations used (turbulence models), which are, in themselves, the cause of additional errors.
- The Boussinesq approximation lies in the belief that the Reynolds Stress tensor behaves in a similar fashion as the Newtonian viscous stress tensor.
- The EVM models assume a linear behavior of the Reynolds stress tensor.
- The main deficiencies of the EVM models are:
 - The assumption of isotropy in shear flows (which is not strictly true),

$$\overline{u'^2} = \overline{v'^2} = \overline{w'^2} = \frac{2}{3}k$$

- The possibility of predicting negative normal shear stresses.

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- Let us recall the Boussinesq approximation.
- Remember, this approximation is the core of all eddy viscosity models (EVM).

$$\tau^R = -\rho \overline{(\mathbf{u}'\mathbf{u}')} = 2\mu_t \bar{\mathbf{S}}^R - \frac{2}{3}\rho k \mathbf{I} = \mu_t [\nabla \bar{\mathbf{u}} + \nabla \bar{\mathbf{u}}^T] - \frac{2}{3}\rho k \mathbf{I}$$

Reynolds averaged strain-rate tensor

Identity matrix (or Kronecker delta).

$$\bar{\mathbf{S}}^R = \frac{1}{2} [\nabla \bar{\mathbf{u}} + \nabla \bar{\mathbf{u}}^T] \quad k = \frac{1}{2} \overline{\mathbf{u}' \cdot \mathbf{u}'} = \frac{1}{2} (\overline{u'^2} + \overline{v'^2} + \overline{w'^2})$$

$$\mathbf{I} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

Which is equivalent to the Kronecker delta

$$\delta_{ij} = \begin{cases} 1, & \text{if } i = j, \\ 0, & \text{if } i \neq j. \end{cases}$$

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- Let us summarize a few applications where the Boussinesq approximation is not very accurate,
 - Poor performance in flows with large extra strains, *e.g.*, curved surfaces, strong vorticity, swirling flows.
 - Rotating flows, *e.g.*, turbomachinery, wind turbines.
 - Impinging flows.
 - Highly anisotropic flows and flows with secondary motions, *e.g.*, fully developed flows in non-circular ducts.
 - Non-local equilibrium and flow separation, *e.g.*, airfoil in stall, dynamic stall.
 - Complex three-dimensional flows.
 - Residual turbulent viscosity near the walls.
- Many EVM models has been developed and improved along the years, so they address the shortcomings of the Boussinesq approximation.

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- In spite of the theoretical weakness of the Boussinesq approximation, it does produce reasonable results for a large number of flows.
- EVM models are the cornerstone of turbulence modeling in industrial applications.
- EVM is an area of active research and new ideas and palliatives to the known deficiencies continues to emerge.
- The deficiencies are found when comparing the numerical results with experimental results, by identifying abnormal behaviors when analyzing the budgets of the transport equations of the turbulence quantities and the shear stresses budget, or by noticing deviations of the non-dimensional velocities from the empirical correlations (law of the wall), among many.
- Many of the corrections take the form of:
 - Additional source terms.
 - Extra terms in the transport equations.
 - Corrective factors (damping, blending, limiting) in some of the terms of the transport equations.

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- Let us briefly overview a few of the remedies to some of the problems found with EVM.
- Remember, these corrections work by adding source terms, extra terms, or corrective factors (damping, blending, limiting) in some of the terms of the transport equations.
- By looking at the general form of the transport equations of the turbulent quantities, any term can be corrected.

$$\underbrace{\nabla_t \phi}_{\text{Transient term}} + \underbrace{\nabla \cdot \bar{\mathbf{u}} \phi}_{\text{Convection}} = \underbrace{P^\phi}_{\text{Production}} + \underbrace{\epsilon^\phi}_{\text{Dissipation}} + \underbrace{D^\phi}_{\text{Diffusion}} + \underbrace{S^\phi}_{\text{Source terms}}$$

- Finally, we will briefly outline some corrective methods, we are not going to elaborate on the closure coefficients or the closure relationships.
- If you are interested in getting more details about one of the methods, we invite you to read the references cited.

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Production limiters

- A disadvantage of two-equation turbulence models is the excessive generation of the turbulent kinetic energy in the vicinity of stagnation points.
- Many corrections have been proposed to avoid this problem.
- Let us discuss a production limiter approach originally proposed by Menter [1].
- In order to avoid the buildup of turbulent kinetic energy in the stagnation regions, the production term in the turbulence equations can be limited as follows,

$$P_k = \min (P_k, C_{lim} \rho \epsilon)$$

- Where the coefficient C_{lim} has a default value of 10.
- This limiter does not affect the shear layer performance of the model, it only avoids the buildup of TKE in stagnation points in aerodynamic simulations.
- Another formulation is based on the work of Kato and Launder [2].

[1] F. R. Menter. Two-Equation Eddy-Viscosity Turbulence Models for Engineering Applications. AIAA Journal. 32(8). 1598–1605. August 1994.

[2] M. Kato and B. E. Launder. The modelling of turbulent flow around stationary and vibrating square cylinders. Ninth Symposium on Turbulent Shear Flows. Kyoto, Japan. August 16-18, 1993.

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Curvature correction

- One drawback of the EVM is that they are insensitive to streamline curvature and system rotation, which play a significant role in many turbulent flows of practical interest.
- Spalart and Shur [1], Shur et al. [2], and Smirnov and Menter [3], have proposed a modification to the production term in order to sensitize EVM models to the effects of streamline curvature and system rotation,

$$P_k \rightarrow P_k f_r$$

$$f_r = \max \left[0, 1 + C_{curv} \left(\tilde{f}_r - 1 \right) \right] \quad \tilde{f}_r = \max \left[\min (f_{rotation}, 1.25), 0 \right]$$

$$f_{rotation} = (1 + C_{r1}) \frac{2r^*}{1 + r^*} \left[1 - C_{r3} \tan^{-1} (C_{r2} \tilde{r}) \right] - C_{r1}$$

C_{curv}

Coefficient to allow influence the strength of the curvature correction. This is the coefficient that controls this correction.

- The rest of the closure coefficient and relationships can be found in reference [1].

[1] P. R. Spalart, M. L. Shur. On the Sensitization of Turbulence Models to Rotation and Curvature. Aerospace Sci. Tech. 1(5). 297–302. 1997.

[2] M. L. Shur, M. K. Strelets, A. K. Travin, P. R. Spalart. Turbulence Modeling in Rotating and Curved Channels: Assessing the Spalart-Shur Correction. AIAA Journal. 38(5). 2000.

[3] P. E. Smirnov, F. R. Menter. Sensitization of the SST Turbulence Model to Rotation and Curvature by Applying the Spalart-Shur Correction Term. ASME Paper GT 2008-50480. 2008.

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Prediction/correction of excessive heat transfer

- In separated or impinging flows, the near-wall length-scale can become too large, resulting in excessively high levels of near-wall turbulence.
- To remedy this behavior, Yap [1] introduced an extra source term into the dissipation rate equation.
- The modification, which is used on the ϵ transport equation is given as follows,

$$S_{\epsilon} = 0.83 \frac{\epsilon^2}{k} \max \left[\left(\frac{l}{\epsilon l_e} - 1 \right) \left(\frac{l}{\epsilon l_e} \right)^2, 0 \right] \quad \text{where} \quad l = \frac{k^{1.5}}{\epsilon} \quad l_e = C_{\mu}^{-0.75} \kappa y$$

- This correction is useful to prevent excessive heat transfer at re-attachment and stagnation points.
- To eliminate the dependence of the above source term on the wall distance, a differential form of the length-scale correction was proposed in reference [2].

[1] C. Yap. Turbulent heat and momentum transfer in recirculating and impinging flows . Ph.D. Thesis Manchester University. 1987.

[2] H. Iacovides, M. Raisee. Recent progress in the computation of flow and heat transfer in internal cooling passages of gas turbine blades. Int. J. Heat Fluid Flow 20:320–328. 1999.

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Realizability

- The term realizability is related to physically tenable Reynolds stress predictions.
- In easier terms, a realizable stated is one in which none of the normal stresses becomes negative, *i.e.*,

$$\overline{u'^2}, \overline{v'^2}, \overline{w'^2} \geq 0$$

- In addition, realizability also satisfy the following condition,

$$\frac{\overline{u'v'}}{\sqrt{\overline{u'^2}}\sqrt{\overline{v'^2}}}, \frac{\overline{u'w'}}{\sqrt{\overline{u'^2}}\sqrt{\overline{w'^2}}}, \frac{\overline{v'w'}}{\sqrt{\overline{v'^2}}\sqrt{\overline{w'^2}}} < 1$$

- The previous relations are referred to as the Schwartz inequality.
- These inequalities can be enforced in EVM models.
- The most known model that satisfy this inequality is the Realizable $k - \epsilon$.

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Damping functions near the walls

- Near the walls, in the viscous sublayer, the turbulent viscosity should exponentially damp to zero.
- Some EVM models might have prediction capability issues since they predict non-zero values (or large values) of turbulent viscosity near the walls.
- To correct this behavior, many EVM use damping functions near the walls.
- For the example, the Van Driest damping function commonly used in mixing length models and some LES models, reads as follows,

$$l = \kappa y \left[1 - e^{-y^+/A^+} \right]$$

$$A^+ = 26$$

This coefficient depends on the pressure gradient

l is use as the length scale to compute the turbulent viscosity

- Another damping function used in the low-RE $k - \epsilon$ turbulence model by Jones and Launder [1], is written as follows,

$$\nu_t = C_\mu f_\mu \frac{k^2}{\epsilon}$$

$$f_\mu = e^{-2.5/(1+Re_T/50)}$$

$$Re_T = \frac{k^2}{\nu \epsilon}$$