

SUPPORTING INFORMATION FOR

# Substrate Chemistry and Lattice Effects in Vapor Transport Growth of Vanadium Dioxide Microcrystals

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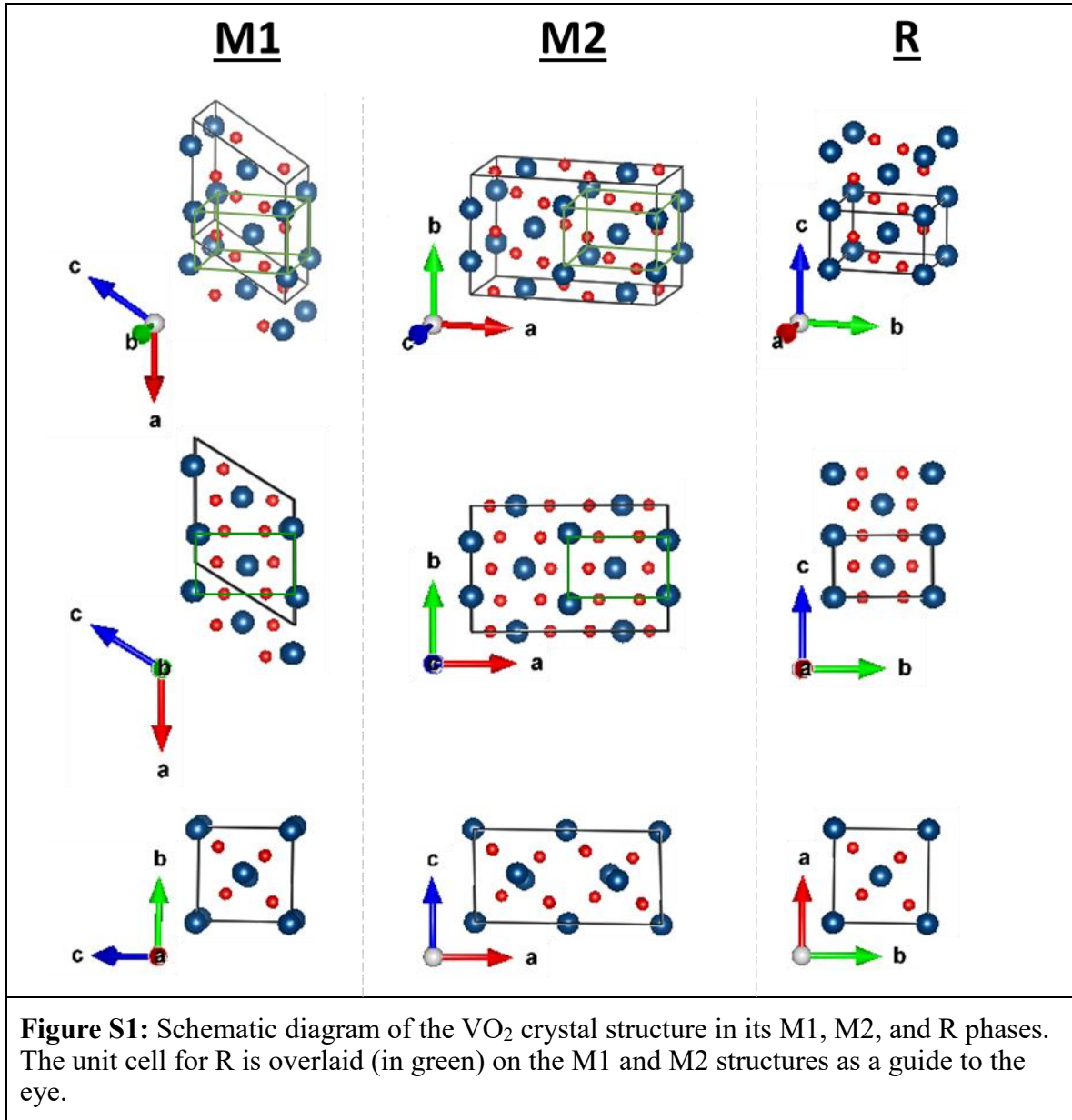
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## S1: VO<sub>2</sub> CRYSTAL LATTICE AND PHASES

Figure S1 compares the unit cells of the M1, M2, and R phases of VO<sub>2</sub>. The corresponding lattice parameters referenced throughout this work are presenting in Table S1.



**Table S1:** Lattice parameters for the M1, M2, and R VO<sub>2</sub> phases

Phase	Space Group	a (Å)	b (Å)	c (Å)	$\alpha(^{\circ})$	$\beta(^{\circ})$	$\gamma(^{\circ})$	Ref.
M1	P 21/c	5.7517	4.5378	5.3825	90	122.646	90	<sup>1</sup>
M2	C 2/m	9.0600	5.8000	4.5217	90	91.850	90	<sup>2</sup>
R	P 42/m n m	4.5546	4.5546	2.8514	90	90	90	<sup>3</sup>

## S2: REVIEW OF PROGRESS IN VAPOR TRANSPORT GROWTH

The optimization of different parameters involved in vapor transport growth has progressed significantly. This method was first developed using  $\text{VO}_2$  as a precursor<sup>4-11</sup>, but  $\text{V}_2\text{O}_5$  has become the standard precursor powder because of its lower melting point.  $\text{VO}_2$  melts at  $1967^\circ\text{C}$ , so using  $\text{V}_2\text{O}_5$ , which melts at about  $690^\circ\text{C}$ , allows for a larger flux of vapor making for a more efficient growth<sup>12-13</sup>. Results of previous growths also show that varying growth conditions affects the morphology of the products. Higher temperatures tend to yield denser nanowires which may show directional growth or even larger nanoplatelets, whereas at lower temperatures, fewer, more loosely attached nanowires tend to form<sup>8, 14</sup>. In at least one study, use of a two-step method, wherein oxygen flow is introduced after the initial temperature ramp, yielded a higher density of nanowires as well<sup>7</sup>.

Choice of substrate has been extensively explored as a parameter of crystal growth. Several studies have focused on the effects of lattice match, showing that varying the lattice structure of the substrate may alter the shape and orientation of  $\text{VO}_2$  crystals. Table S2 summarizes substrates that have been used and how lattice match affects the crystals they yield. Generally,  $\text{VO}_2$  crystals tend to grow along the rutile c-axis  $[001]_R$  with facets along the low-surface-energy  $\{110\}_R$  family of planes<sup>11</sup>.

The most commonly used substrates for  $\text{VO}_2$  crystal growth are silicon and silica. Vanadium dioxide crystals grown on pure  $\text{SiO}_2$  tend to grow in specific directions depending on the cut used (with  $[001]_R$  aligned to  $\langle 2\bar{1}10 \rangle_{\text{SiO}_2}$  on z-cut  $\text{SiO}_2$ <sup>14</sup> and to  $\langle 1\bar{1}00 \rangle_{\text{SiO}_2}$  on x-cut<sup>12</sup>). Silicon oxide on Si, pure Si, and  $\text{Si}_3\text{N}_4$  have produced crystals with unique structural characteristics, but with random orientation in the plane of the substrate<sup>4-5, 15-17</sup>. Growth on  $\text{TiO}_2$  (100) and  $\text{TiO}_2$  (110) also results in crystals with different structural characteristics, having unique cross-sectional shapes, unusual crystal facets parallel to the substrate, and specific in-plane orientation<sup>11</sup>. Another widely-used substrate is  $\text{Al}_2\text{O}_3$ . In addition to varying structural characteristics, crystals grown on  $\text{Al}_2\text{O}_3$  show definite evidence of preferred growth direction relative to the substrate, depending on the cut used<sup>6, 10</sup>. On c-cut, crystals grow at  $60^\circ$  angles relative to each other<sup>6</sup>, whereas in r-cut and a-cut, crystals grow directionally out of plane relative to the substrate<sup>10</sup>. Across different crystalline substrates, the arrangement of atoms on the surface affects the shape and morphology of the crystals on a micron-scale.

Other substrate-dependent effects during crystal growth play a significant role in determining the properties of the resulting  $\text{VO}_2$  crystals. In particular, chemical interactions between the substrates and vanadium precursors affect the growth mechanism. On substrates with  $\text{SiO}_2$  at the surface, a eutectic region forms between newly formed  $\text{VO}_x$  droplets and  $\text{VO}_2$  wires which dissolves a small amount of surrounding material, embedding the wire into the substrate. This facilitates unobstructed growth which leads to larger crystals<sup>4, 17</sup>, and firmly clamps the crystals to the substrate. On  $\text{TiO}_2$  substrates, the precursor vapor wets the surface during growth instead of forming droplets<sup>17</sup>. These observations show that interfacial interactions between liquid vanadium oxides and the substrate surface changes the growth mechanism on these substrates, which may result in different crystal morphologies. In addition, some crystals

grown on  $\text{Al}_2\text{O}_3$  and  $\text{SiO}_2$  have been shown to exist in the M2 phase at room temperature<sup>7, 9, 18</sup>. Although this has often been attributed to substrate-induced strain, the M2 phase persists in crystals removed from  $\text{Al}_2\text{O}_3$  substrates (and thus relieved of substrate-induced strain). Moreover, the M2 phase has also been stabilized by aluminum doping<sup>19</sup>, which suggests the possibility of a chemical effect.

<b>Table S2:</b> Summary of substrate effects on $\text{VO}_2$ crystal growth		
<b>Substrate</b>	<b>Evidence of Lattice Match / Orientation</b>	<b>Evidence of Chemistry Effects</b>
$\text{SiO}_2$ z-cut (0001) <sub><math>\text{SiO}_2</math></sub>	$[001]_R \parallel \langle 2\bar{1}\bar{1}0 \rangle_{\text{SiO}_2}$ <sup>14</sup> $\{110\}_R \parallel \{0001\}_{\text{SiO}_2}$	Molten $\text{V}_2\text{O}_5$ forms eutectic with $\text{SiO}_2$ ; eutectic layer facilitates growth; resulting $\text{VO}_2$ crystals embedded in $\text{SiO}_2$ <sup>4, 17</sup>
$\text{SiO}_2$ x-cut (11 $\bar{2}$ 0) <sub><math>\text{SiO}_2</math></sub>	$[001]_R \parallel \langle 1\bar{1}00 \rangle_{\text{SiO}_2}$ <sup>12</sup> $\{110\}_R \parallel \{11\bar{2}0\}_{\text{SiO}_2}$	
$\text{SiO}_2$ (amorphous) on Si	$\{110\}_R$ <sup>4, 17</sup> and $\{011\}_R$ <sup>4</sup> parallel to substrate; square-cross section <sup>4, 16</sup> ; random orientation in plane <sup>17</sup>	
Si (w/ native oxide)	Random orientation in plane <sup>6</sup>	Si-assisted reduction: $2\text{V}_2\text{O}_5(l) + \text{Si}(s) \rightarrow 4\text{VO}_2(s) + \text{SiO}_2$ Leads to small, dense nuclei <sup>17</sup>
$\text{Si}_3\text{N}_4$	Rectangular cross section; $\{110\}_R$ parallel to substrate <sup>5</sup>	Higher density of nanowires compared to $\text{SiO}_2$ substrates <sup>5</sup>
GaN (0001) <sub>GaN</sub>	In-plane orientation reflects hexagonal substrate symmetry <sup>17</sup>	Molten vanadia wets GaN surface <sup>17</sup>
$\text{TiO}_2$ (110) <sub><math>\text{TiO}_2</math></sub>	Triangular cross section; exposed facets are $\{100\}_R$ planes; $[001]_R \parallel \langle 001 \rangle_{\text{TiO}_2}$ <sup>11</sup> $\{110\}_R \parallel \{110\}_{\text{TiO}_2}$	Molten vanadia wets $\text{TiO}_2$ surface <sup>17</sup>
$\text{TiO}_2$ (100) <sub><math>\text{TiO}_2</math></sub>	Rectangular cross section; exposed facets are $\{100\}_R$ planes; $[001]_R \parallel \langle 001 \rangle_{\text{TiO}_2}$ <sup>11</sup> $\{100\}_R \parallel \{100\}_{\text{TiO}_2}$	
$\text{Al}_2\text{O}_3$ c-cut (0001) <sub><math>\text{Al}_2\text{O}_3</math></sub>	Triangular cross-section <sup>10-11</sup> , $[001]_R \parallel \langle 11\bar{2}0 \rangle_{\text{Al}_2\text{O}_3}$ <sup>10</sup> , $\{100\}_R \parallel \{0001\}_{\text{Al}_2\text{O}_3}$	M2 phase stabilized at room temperature, attributed to strain effects <sup>7, 9</sup>
$\text{Al}_2\text{O}_3$ r-cut (11 $\bar{2}$ 0) <sub><math>\text{Al}_2\text{O}_3</math></sub>	Square-cross section <sup>10</sup> ; growth axis out of plane, 60° relative to substrate surface, unique in-plane orientation <sup>7, 10</sup>	
$\text{Al}_2\text{O}_3$ a-cut (11 $\bar{2}$ 0) <sub><math>\text{Al}_2\text{O}_3</math></sub>	Square-cross section; growth axis out of plane, 3-fold in-plane orientation symmetry <sup>10</sup>	
$\text{Al}_2\text{O}_3$ m-cut (10 $\bar{1}$ 0) <sub><math>\text{Al}_2\text{O}_3</math></sub>	Nanoblocks and nanowires appear to exhibit preferred orientation <sup>20</sup>	
$\text{Al}_2\text{O}_3$ (amorphous) on $\text{SiO}_2$		Al-doping stabilizes M2-phase at room temperature <sup>19</sup>

### S3: SAMPLE AND PRECURSOR LOADING

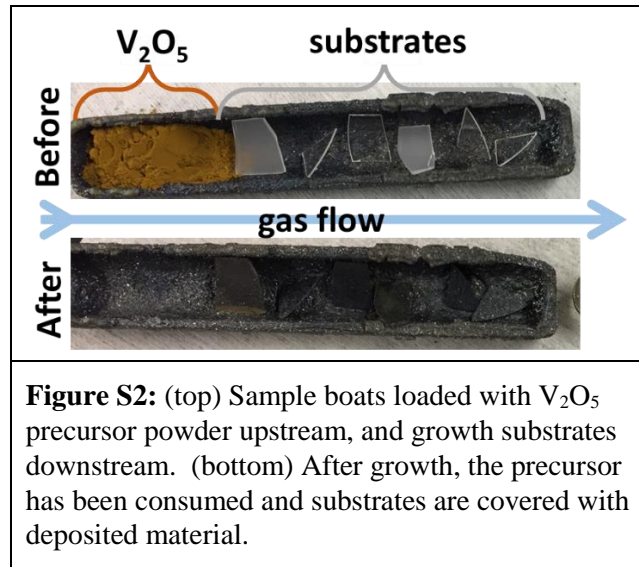
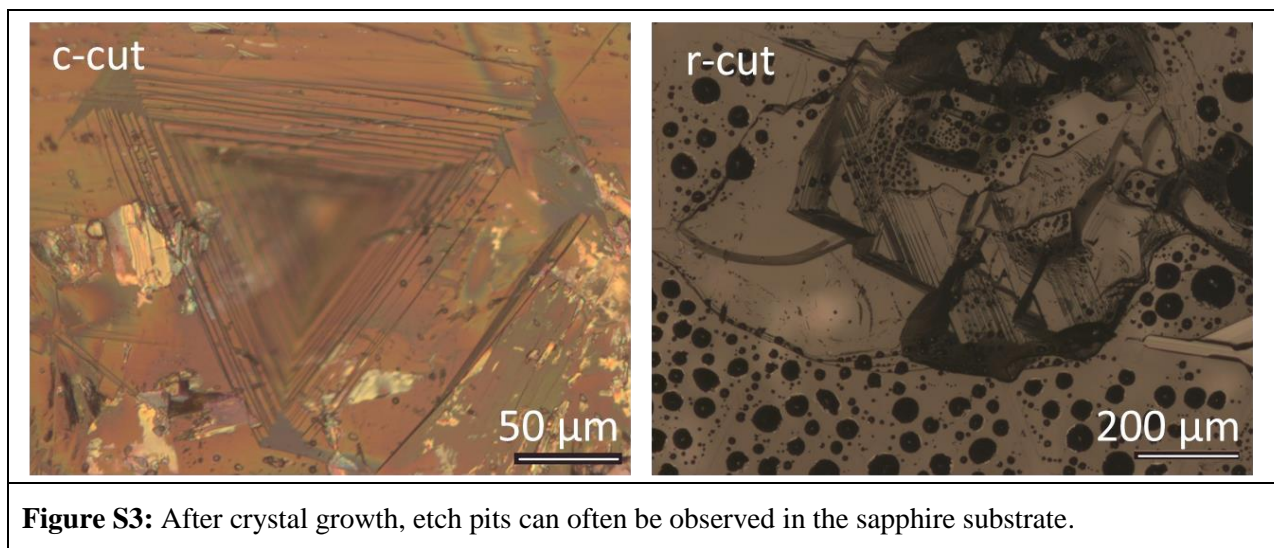


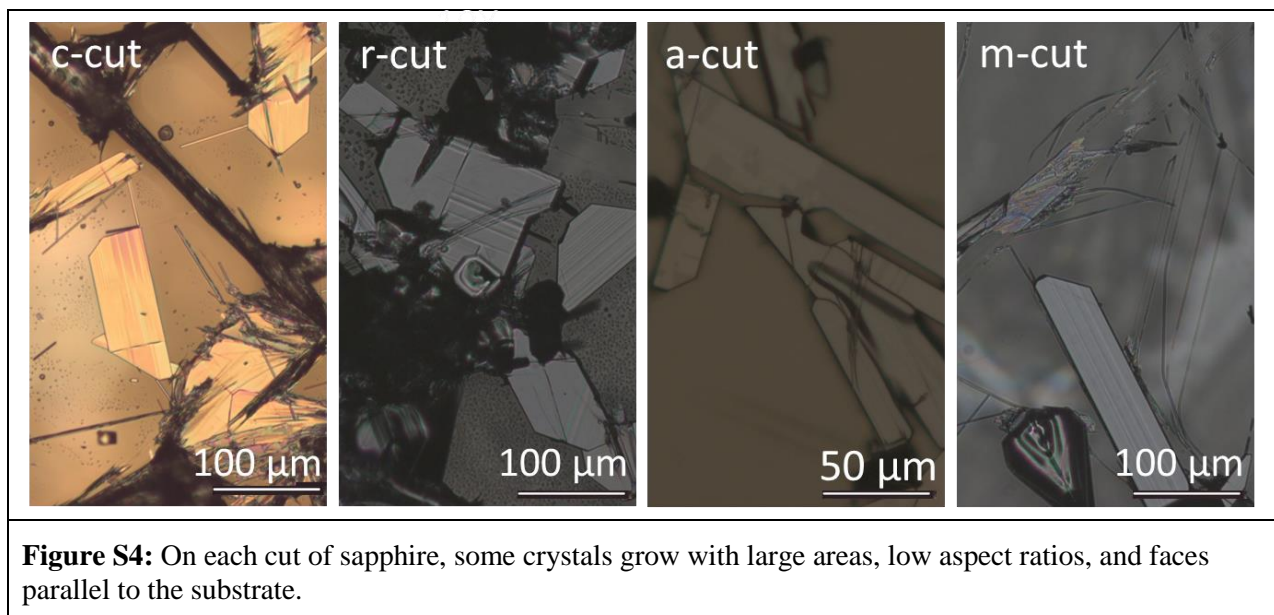
Figure S2 shows an example of the samples before and after the growth process. Orange vanadium pentoxide ( $V_2O_5$ ) powder is loaded into one end of the boat, and growth substrates into the other. During the growth process,  $V_2O_5$  vapor is carried to the substrates by a directional gas flow. After growth, the  $V_2O_5$  is consumed and the substrates are darkened with a coating of  $VO_2$  crystals.

#### S4: ETCH PITS IN SAPPHIRE



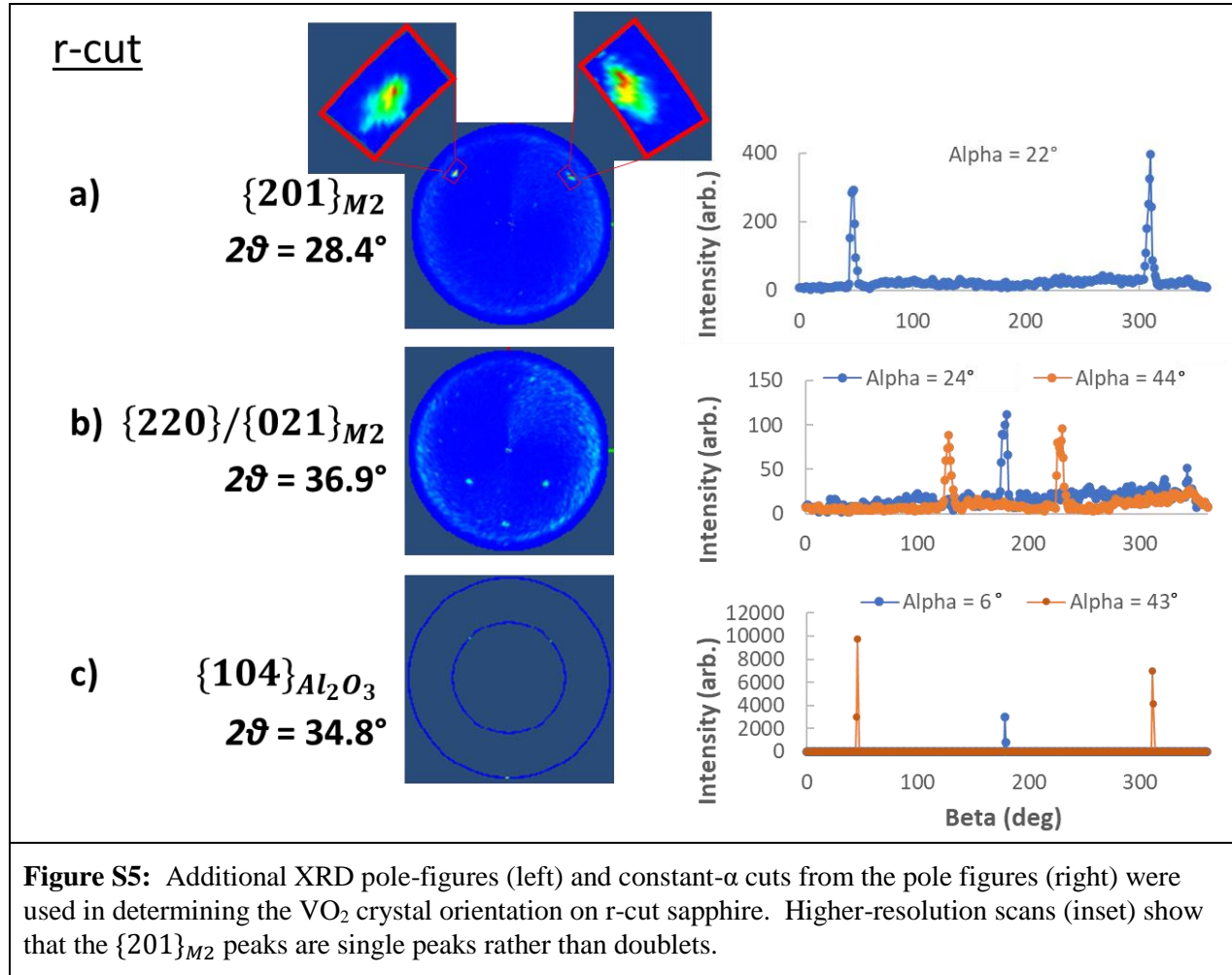
After the growth process, pits can be observed in the sapphire substrate, examples of which are shown in Figure S3. These have a similar appearance to etch pits in sapphire<sup>21</sup>, with a tendency to have stepped edges and to be shaped differently on different cuts, reflecting the difference in etch rates on different lattice planes. These pits are not observed on untreated substrates, or on substrates heated in the absence of  $V_2O_5$  precursor. We attribute them to the corrosive action of molten  $V_2O_5$ , which has been used to etch  $Al_2O_3$  surfaces<sup>22-23</sup>.

S5: EXAMPLES OF LARGE, FLAT, LOW-ASPECT-RATIO MICROCRYSTALS ON SAPPHIRE



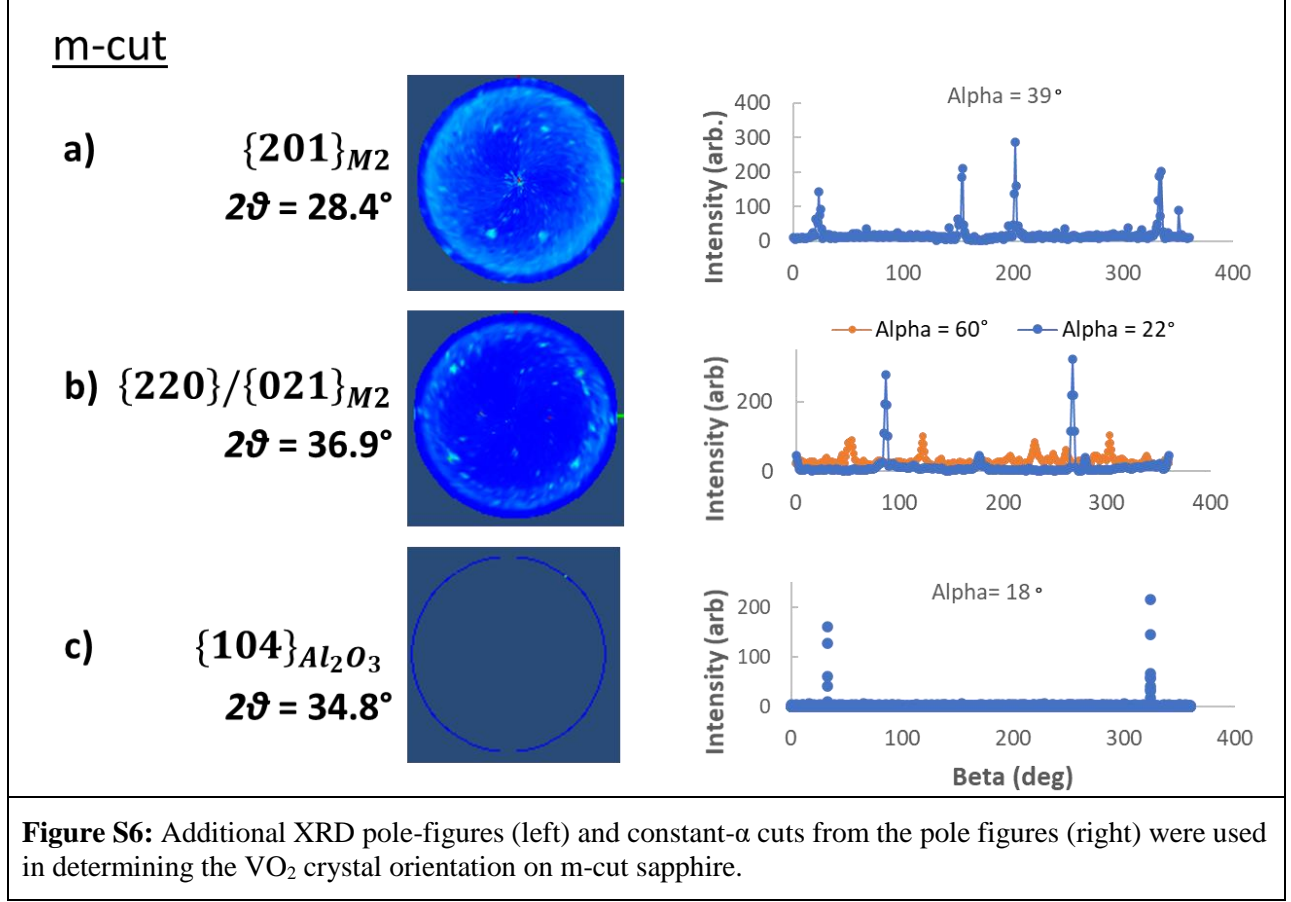
As discussed in the main text, both highly-oriented nanowires and un-oriented micro-platelets can be observed on each cut of sapphire. Examples of the latter are shown in Figure S4.



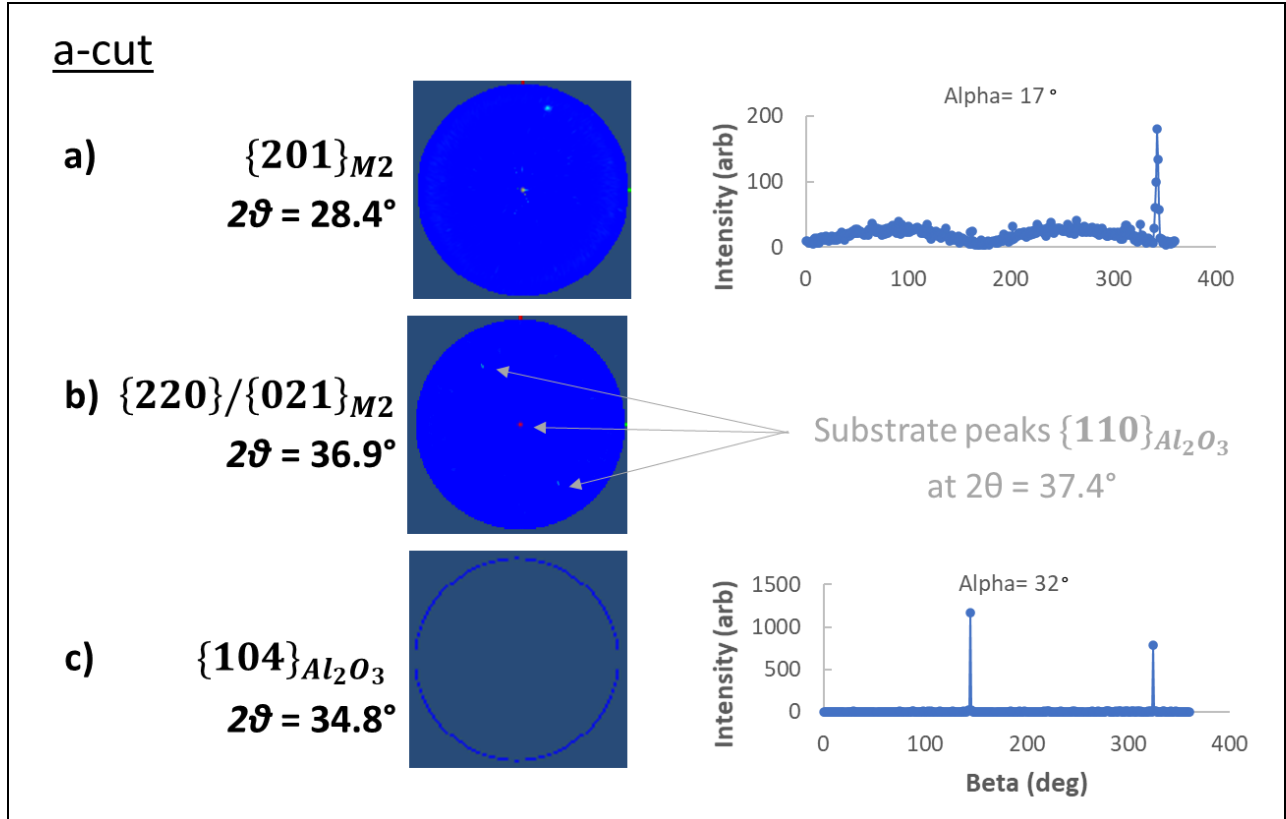


On r-cut  $Al_2O_3$ , the  $\{201\}_{M2} \equiv \{110\}_R$  family of planes shows one peak at  $\alpha=90^\circ$  ( $\{110\}_R$  parallel to substrate) and two at  $\alpha=22^\circ$  with  $\beta=-2^\circ \pm 48^\circ$ . The  $\{220\}/\{021\}_{M2} \equiv \{011\}_R$  planes show four clear peaks:  $\alpha=90^\circ$ ,  $\alpha=45^\circ$  with  $\beta=-1^\circ \pm 129^\circ$ , and  $\alpha=24^\circ$  with  $\beta=180^\circ$ . These four peaks correspond to the two  $\{110\}_R$  peaks at  $\alpha=22^\circ$ . On the other hand, no  $\{011\}_R$  peaks are observed that correspond to the  $\{110\}_R$  peak at  $\alpha=90^\circ$ , which implies that these crystals have no in-plane orientation. Thus, some of the  $VO_2$  crystals on r-cut  $Al_2O_3$  have  $\{110\}_R$  parallel to the substrate with no in-plane orientation, while others have  $\{011\}_R$  parallel to the substrate with one unique in-plane orientation. To establish the substrate orientation,  $\{104\}_{Al_2O_3}$  peaks are observed at  $\alpha=43^\circ$  with  $\beta=-2^\circ \pm 48^\circ$  and  $\alpha=6^\circ$  with  $\beta=178^\circ$ . This establishes the orientation of the  $VO_2$  crystals relative to the substrate, as listed in Table 2 of the main text.

Notably, this orientation is not the one that would yield the smallest lattice mismatch theoretically. If instead the orientation was  $\langle 11\bar{1} \rangle_R \parallel \langle 0\bar{2}2\bar{1} \rangle_{Al_2O_3}$ , the calculated mismatch would be only 0.3% (1.4%). This orientation is very close to that we observe, having only a  $\pm 2.5^\circ$  shift in  $\beta$ ; however, this would lead to a  $5^\circ$  peak splitting, which is not observed in our data, even with higher-resolution measurements (Figure S5a, insets).



On m-cut  $Al_2O_3$ , the  $\{201\}_{M2} \equiv \{110\}_R$  family of planes shows one peak at  $\alpha=90^\circ$  ( $\{110\}_R$  parallel to substrate) and four at  $\alpha=39^\circ$  with  $\beta=\pm 24^\circ$  or  $\beta=\pm 155^\circ$ . The  $\{220\}/\{021\}_{M2} \equiv \{011\}_R$  planes show four clear peaks, all at  $\alpha=21^\circ$  with  $\beta=1^\circ \pm 54^\circ$  or  $\beta=1^\circ \pm 125^\circ$ . These four peaks correspond to the two  $\{110\}_R$  peaks at  $\alpha=39^\circ$ , and arise from  $VO_2$  crystals in two symmetry-equivalent orientations, which can be fully identified by comparison to the substrate  $\{104\}_{Al_2O_3}$  peaks at  $\alpha=18^\circ$  and  $\beta=-2^\circ \pm 34^\circ$  (see main text, Table 2). Once again, no  $\{011\}_R$  peaks can be identified to correspond to the  $\{110\}_R$  peaks at  $\alpha=90^\circ$ , indicating that the crystals with  $\{110\}_R$  parallel to the substrate have no definite in-plane orientation.

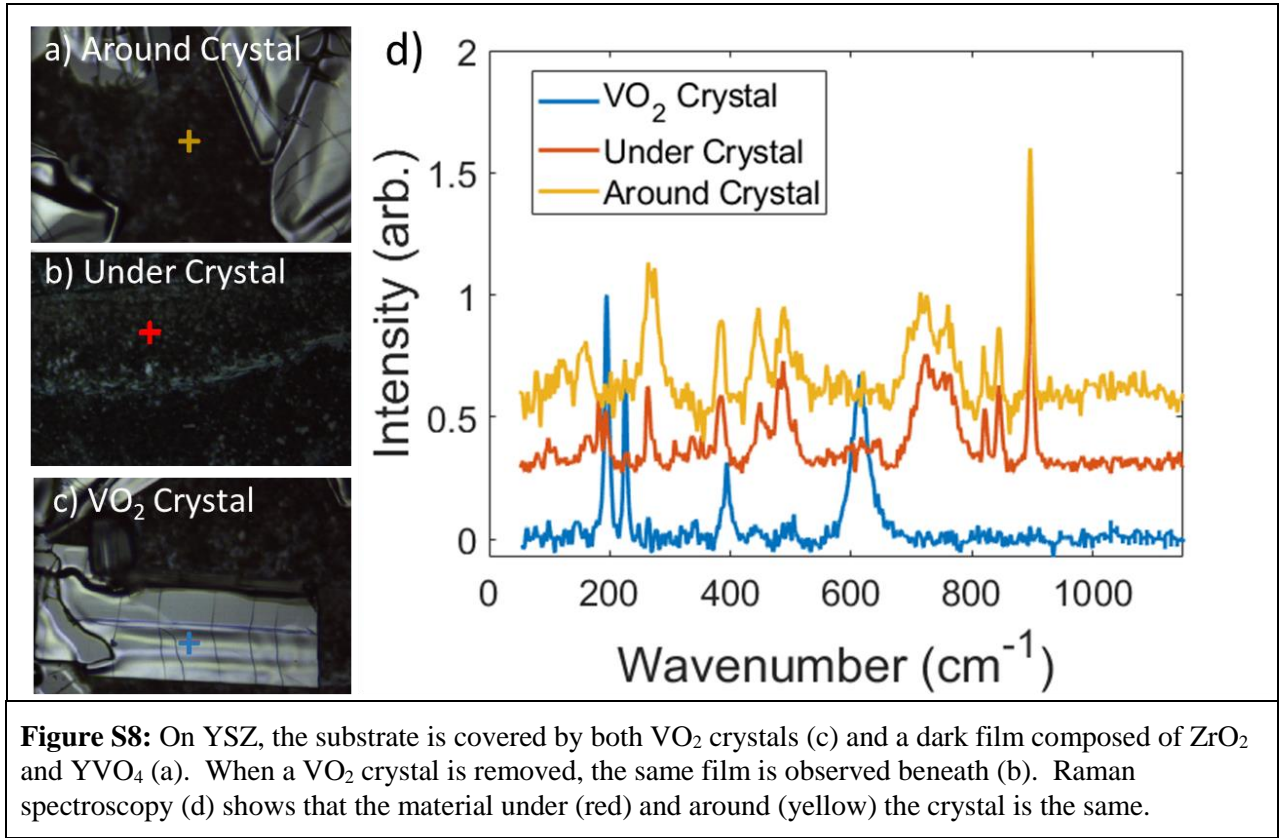


**Figure S7:** Additional XRD pole-figures (left) and constant- $\alpha$  cuts from the pole figures (right) were used in determining the  $VO_2$  crystal orientation on m-cut sapphire. The peaks that appear in the  $\{220\}/\{021\}_{M2}$  pole figure (at  $2\theta = 36.9^\circ$ ) are positioned where  $\{110\}_{Al_2O_3}$  peaks (at  $2\theta = 37.4^\circ$ ) would be expected, and are likely the result of bleed-over from those reflections.

On a-cut  $Al_2O_3$ , the  $\{201\}_{M2} \equiv \{110\}_R$  family of planes shows, as usual, one peak at  $\alpha=90^\circ$  ( $\{110\}_R$  parallel to substrate), but only one other peak, at  $\alpha=17^\circ$  with  $\beta=-17^\circ$ . Given the substrate symmetry, at least two peaks at  $\alpha=17^\circ$  would be expected for a sufficiently large sample of crystals; however, this sample was more sparsely covered with  $VO_2$  than the others, and it may be that the corresponding symmetric peak was simply too weak to be observed. Similarly, no peaks are observed due to the  $\{220\}/\{021\}_{M2} \equiv \{011\}_R$  family of planes. The peaks that do appear in the pole figure (Figure S6b) can be assigned to bleed-over from the (much more intense)  $\{110\}_{Al_2O_3}$  peaks which have a similar  $2\theta$  value. All of the  $\{011\}_R$  peaks, being intrinsically less intense than the  $\{110\}_R$ , are simply too weak to be observed. However, to fully determine the orientation which produced the one  $\{110\}_R$  peak at  $\alpha=17^\circ$  requires additional information. Standard  $\theta$ - $2\theta$  measurements reveal all the  $VO_2$  planes which are parallel to the substrate surface. Table S3 lists all the M2- $VO_2$  planes identified in  $\theta$ - $2\theta$  measurements on a-cut  $Al_2O_3$ . Some of these may correspond to crystals with no in-plane orientation (as appears to be typical of the  $\{201\}/\{20\bar{1}\}_{M2}$  planes), or to loose, randomly oriented crystals, but we can identify one that is consistent with  $\{110\}_R$  peak at  $\alpha=17^\circ$ , which makes it possible to identify that orientation.

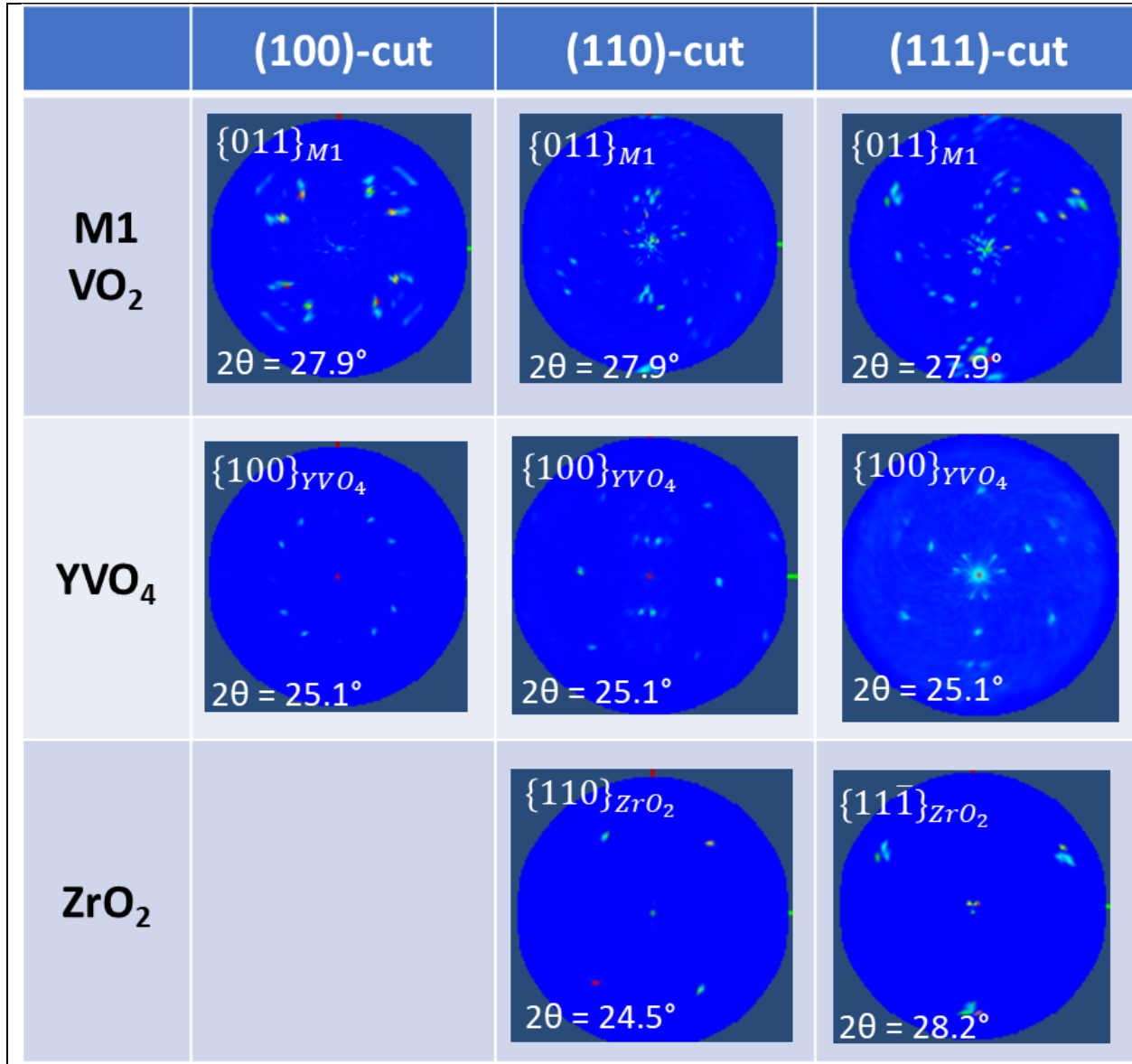
<b>Table S3:</b> M2 VO <sub>2</sub> peaks observed in $\theta$ - $2\theta$ measurements on a-cut sapphire
$\{201\}/\{20\bar{1}\}$
$\{220\}/\{021\}$
$\{400\}/\{002\}$
$\{401\}/\{202\}$

## S7: EVIDENCE THAT A YVO<sub>4</sub> LAYER UNDERLIES VO<sub>2</sub> CRYSTALS ON YSZ



As discussed in the main text, YSZ substrates are covered in a dark, granular film after the crystal growth process. Raman spectroscopy showed this film to be composed of ZrO<sub>2</sub> and YVO<sub>4</sub>. To determine whether this film covers the entire substrate or only that portion not occupied by VO<sub>2</sub> crystals, we remove a VO<sub>2</sub> crystal with adhesive tape and examine the material underneath. The crystal (Figure S7c) exhibits the characteristic Raman spectrum of VO<sub>2</sub> (Figure S7d, blue trace); once it has been removed, the area underneath it has the same visual appearance (Figure S7b) and Raman spectrum (Figure S7d, red trace) as the film that appears elsewhere on the exposed substrate (Figure S7a and S7d, red trace). This shows that the ZrO<sub>2</sub>/YVO<sub>4</sub> film covers the entire substrate before VO<sub>2</sub> crystals form above it.

S8: POLE FIGURES FOR  $\text{YVO}_4$ ,  $\text{ZrO}_2$ , ON YSZ



**Figure S9:** Additional pole figures for growths on all cuts of YSZ show evidence of preferred orientation for  $\text{VO}_2$ ,  $\text{YVO}_4$ , and  $\text{ZrO}_2$ .

Figure S9 compares the pole figures for the  $\text{VO}_2$   $\{011\}_{M1} \equiv \{1\bar{1}0\}_R$  family of planes to representative low-index planes belonging to  $\text{YVO}_4$  and  $\text{ZrO}_2$ . Both  $\text{YVO}_4$  and  $\text{ZrO}_2$  exhibit sharp peaks in well-defined patterns, with a high degree of symmetry reflecting the symmetry of the substrate plane: 4-fold on (100)-cut, 2-fold on (110)-cut, and 3-fold on (111)-cut. This shows that the  $\text{YVO}_4$  and  $\text{ZrO}_2$  are highly-oriented, likely due to a heteroepitaxial relationship to the substrate and each other. Notably, the pole-figure patterns become progressively more complex from  $\text{ZrO}_2$  to  $\text{YVO}_4$  to  $\text{VO}_2$ . We hypothesize that atop the YSZ substrate is an yttria-

depleted  $\text{ZrO}_2$  layer, heteroepitaxial to the substrate, but with multiple distinct orientations arising from its decreased symmetry. Above the  $\text{ZrO}_2$  layer, the yttria reacts with  $\text{V}_2\text{O}_5$  to form a  $\text{YVO}_4$  layer, in turn oriented with respect to the  $\text{ZrO}_2$ . Multiple symmetry-allowed orientations lead to multiplicatively more peaks in the pole-figure pattern.  $\text{VO}_2$  crystals grown on the YSZ layer in turn have preferred orientations relative to it. Finally, we note that since the  $\{011\}_{M1}$  and  $\{11\bar{1}\}_{\text{ZrO}_2}$  planes have similar  $2\theta$  values ( $27.9^\circ$  and  $28.2^\circ$ , respectively) is it possible that some bleed-over from the more-intense  $\text{ZrO}_2$  peaks contributes to the  $\text{VO}_2$  pole figures.



## S9: GROWTH RESULTS ON QUARTZ

Quartz is perhaps the most commonly-used and well-characterized substrate for VO<sub>2</sub> crystal growth by vapor-phase transport. As a baseline comparison for our crystals on other substrates, we present a sample grown on z-cut (0001) quartz. On this cut of quartz, crystals may exhibit a preferred orientation with six-fold symmetry<sup>14</sup>, as noted above; and as on all SiO<sub>2</sub> substrates, VO<sub>2</sub> crystals will be embedded slightly in substrate due to formation of a eutectic during growth<sup>17</sup>.

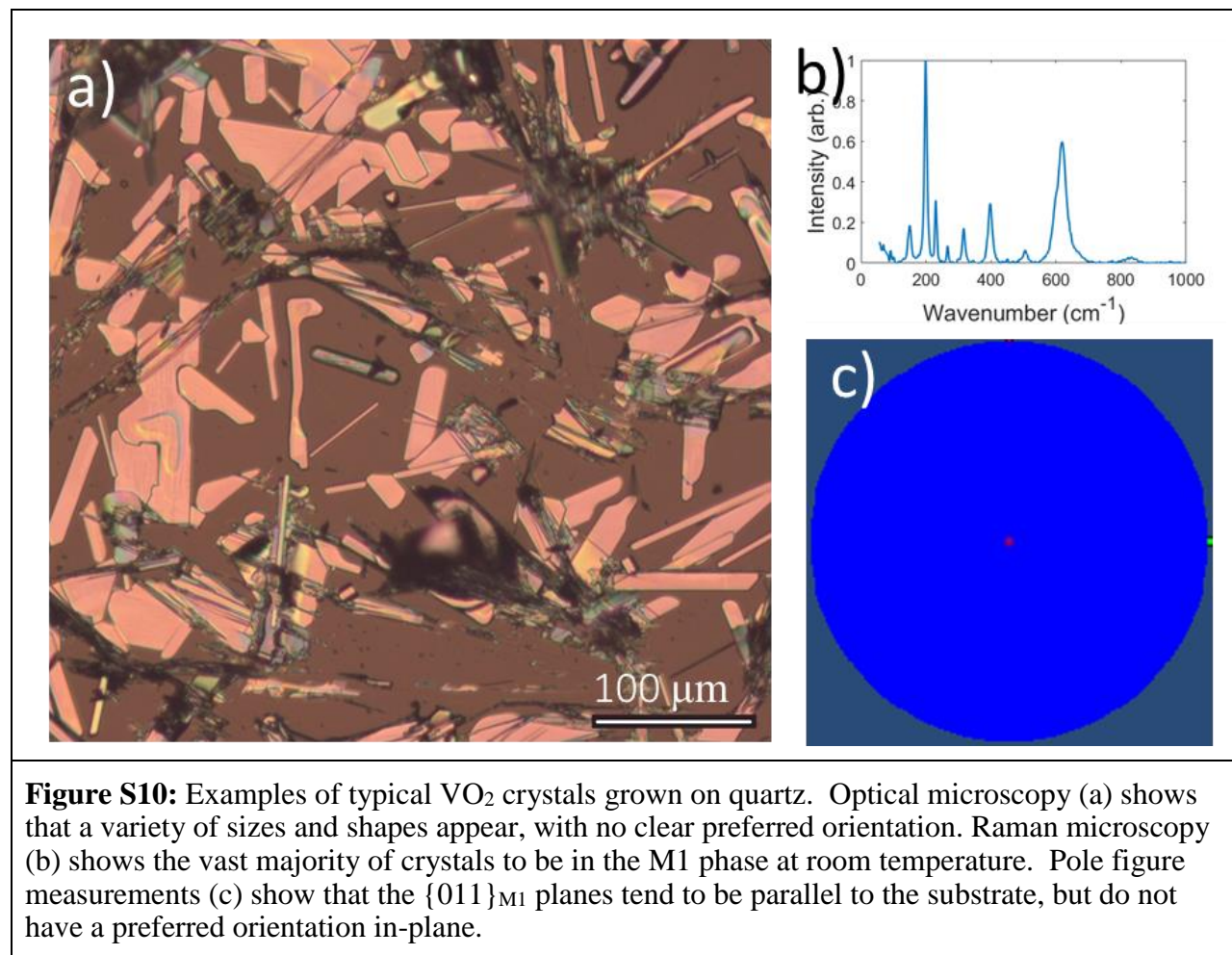
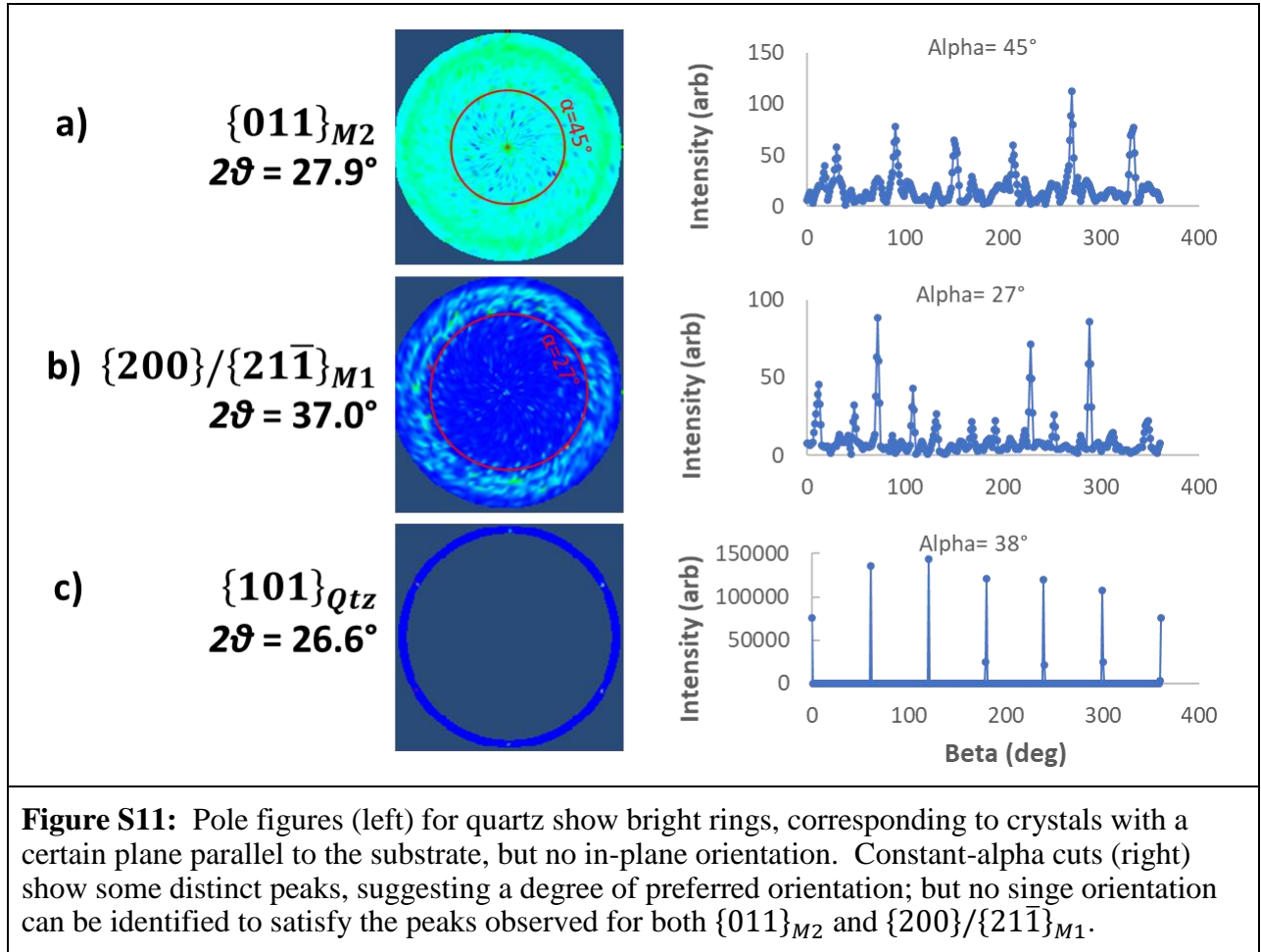


Figure S10 shows example data of our VO<sub>2</sub> crystals grown on z-cut quartz. Polarized optical microscopy (Figure S10a) shows that the crystals lie flat, with upper facets parallel to the substrate surface. The crystals are well adhered to the substrate—due to being embedded into the quartz, as mentioned above—and the substrate-induced strain results in complex ferroelastic twin domains (observed as light and dark mazelike patterns in the microscope image)<sup>24</sup>. There is no obvious preferred orientation, and our crystals are more reminiscent of those others have observed on SiO<sub>2</sub> thermal oxides on Si<sup>4, 17</sup>, than the oriented nanowires observed on z-cut quartz<sup>14</sup>. As we observed in our sapphire growths, it is likely possible for both oriented



nanowires and non-oriented microplatelets to form, especially under different growth conditions. Raman microscopy (Figure S10b) shows that the majority of VO<sub>2</sub> crystals on quartz are in the M1 phase at room temperature, though some are in the T phase due to the influence of substrate strain.

Pole figure measurements (Figure S10c) show that the vast majority of  $\{011\}_{M1} \equiv \{1\bar{1}0\}_R$  planes tend to be parallel to the substrate ( $\alpha=90^\circ$ ). Replotting this figure on a log scale (Figure S11a) reveals an additional bright ring at  $\alpha=45^\circ$  with brighter spots appearing with a six-fold symmetry (reflective of the six-fold symmetry of the z-cut quartz substrate). Another bright ring in this figure at very low  $\alpha$  is background noise present in all pole figures, usually too weak to appear on the scale used. The  $\{200\}/\{21\bar{1}\}_{M1} \equiv \{011\}_R$  show peaks, of variable intensity, at  $\alpha=27^\circ$  with six-fold symmetry:  $\beta=12^\circ+n*60^\circ$  and  $\beta=48^\circ+n*60^\circ$  ( $n=0, 1, 2, 3, 4, 5$ ). However, no VO<sub>2</sub> orientation can be identified to satisfy both  $\{011\}_R$  at  $\alpha=27^\circ$  and  $\{1\bar{1}0\}_R$  at either  $\alpha=90^\circ$  or  $\alpha=45^\circ$ . We conclude that most of our VO<sub>2</sub> crystals on quartz grow with  $\{1\bar{1}0\}_R$  parallel to the substrate with no preferred in-plane orientation, while some crystals grow in at least two distinct orientations. Additional measurements will be required to identify the latter.



## S10: TITANIUM DIOXIDE SUBSTRATE PROPERTIES

For this study, we chose two substrates that exemplify both chemical and lattice-matching effects in a well-understood case ( $\text{Al}_2\text{O}_3$ ) as well as a novel one (YSZ). To similarly analyze all common  $\text{VO}_2$  growth substrates is beyond the scope of this work, but one in particular deserves mention. Though more expensive, titanium dioxide ( $\text{TiO}_2$ ) is commonly regarded as the best substrate for heteroepitaxial  $\text{VO}_2$  growth because it has the same rutile crystal structure with very similar lattice parameters ( $a_{R,\text{TiO}_2} = 4.593\text{\AA}$  and  $c_{R,\text{TiO}_2} = 2.959\text{\AA}$ <sup>25</sup>, compared to  $a_{R,\text{VO}_2} = 4.555\text{\AA}$  and  $c_{R,\text{VO}_2} = 2.851\text{\AA}$ <sup>3</sup>), and has produced highly-oriented films<sup>26-27</sup> and crystals<sup>11</sup>. Table S3 presents predicted lattice mismatch values for epitaxial  $\text{VO}_2$  on several low-index planes of  $\text{TiO}_2$ . Compared to the values calculated for  $\text{VO}_2$  on  $\text{Al}_2\text{O}_3$  (Table 1 of the main text), these are for the most part comparable or better, with one notable exception: the percent lattice match for  $\text{VO}_2$  on c-cut  $\text{Al}_2\text{O}_3$  (according to the orientation we observed) is better than that on  $\{110\}_{\text{TiO}_2}$  or  $\{100\}_{\text{TiO}_2}$  (the two cuts of  $\text{TiO}_2$  used previously for single-crystal  $\text{VO}_2$  growth). Although this represents multi-domain heteroepitaxy on  $\text{Al}_2\text{O}_3$  as opposed to single-domain on  $\text{TiO}_2$ , lattice match on sapphire suffices for our purposes.

<b>Table S4:</b> Predicted orientation and lattice match for $\text{VO}_2$ crystals grown on various cuts of $\text{TiO}_2$				
Substrate Plane	Crystal Orientation	Lattice Mismatch ( $\frac{d_{\text{VO}_2} - d_{\text{TiO}_2}}{d_{\text{TiO}_2}}$ )	$d_{\text{VO}_2}$ (Å)	$d_{\text{TiO}_2}$ (Å)
$\{110\}_R \parallel \{110\}_{\text{TiO}_2}$	$[001]_R \parallel \langle 001 \rangle_{\text{TiO}_2}$	-3.6%	2.851	2.959
	$[\bar{1}10]_R \parallel \langle \bar{1}10 \rangle_{\text{TiO}_2}$	-0.8%	6.442	6.496
$\{100\}_R \parallel \{100\}_{\text{TiO}_2}$	$[001]_R \parallel \langle 001 \rangle_{\text{TiO}_2}$	-3.6%	2.851	2.959
	$[010]_R \parallel \langle 010 \rangle_{\text{TiO}_2}$	-0.8%	4.555	4.593
$\{001\}_R \parallel \{001\}_{\text{TiO}_2}$	$[100]_R \parallel \langle 100 \rangle_{\text{TiO}_2}$	-0.8%	4.555	4.593
	$[010]_R \parallel \langle 010 \rangle_{\text{TiO}_2}$	-0.8%	4.555	4.593
$\{011\}_R \parallel \{011\}_{\text{TiO}_2}$	$[100]_R \parallel \langle 100 \rangle_{\text{TiO}_2}$	-0.8%	4.555	4.593
	$[01\bar{1}]_R \parallel \langle 01\bar{1} \rangle_{\text{TiO}_2}$	-1.6%	5.374	5.464

As with  $\text{Al}_2\text{O}_3$  and YSZ, there is a possibility of chemical reaction between  $\text{TiO}_2$  and  $\text{V}_2\text{O}_5$ . Studies on  $\text{TiO}_2$ -supported  $\text{V}_2\text{O}_5$  catalysts show that at high temperatures vanadia catalyzes the  $\text{TiO}_2$  anatase-to-rutile transition, and  $\text{V}^{4+}$  is incorporated into the rutile  $\text{TiO}_2$  lattice, yielding  $\text{V}_x\text{Ti}_{1-x}\text{O}_2$  at the interface<sup>28-30</sup>. Also,  $\text{V}_2\text{O}_5$  is reported to react with titania-stabilized zirconia ( $\text{TiSZ}$ ) to produce  $\text{TiVO}_4$ <sup>31</sup>. On the other hand, several corrosion studies on Ti-containing thermal barrier coatings suggest that  $\text{TiO}_2$  is less reactive with vanadia than some other metal oxides, including alumina<sup>32-34</sup>. The possibility of chemical reactions occurring at the  $\text{TiO}_2$ - $\text{V}_2\text{O}_5$  interface during crystal growth certainly deserves further study, but sapphire provides a better example of substrate-vanadia reactions for this study.

There are few experimental studies of the  $\text{TiO}_2$  surface energy, but computational studies predict values ranging from 0.3 to 1.8 J/m<sup>2</sup> for the  $\{110\}_{\text{TiO}_2}$  plane, depending on the calculation

method, and higher for other planes <sup>35-38</sup>. The highest of these predicted values are larger than those used in our study, and might lead to good substrate wetting, as observed in one crystal growth study <sup>17</sup>.

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