# **Supporting Information**

## **Molecular Simulation of Tracer and Self-Diffusion in Entangled Polymers**

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# 1 Summary of Systems Investigated

MD	BFM	SS
$50 \times 500$	100 × 1080	5  imes 20
$100 \times 500$	$200 \times 540$	10  imes 20
$200 \times 250$	225  imes 480	20  imes 20
	250  imes 432	35  imes 20
		50  imes 20
		$100 \times 20$
		$150 \times 20$

The systems investigated in this paper are tabulated.

Table 1: Monodisperse Melts for Self-Diffusion Studies: Systems  $(N \times n)$ , where  $N = N_p = N_m$  is the number of beads per chain, and n is the total number of chains in the ensemble.

$N_p \times n_p$	$N_m \times n_m$	BB	СМ
75  imes 49	30 × 1111	X	x
75  imes 49	75  imes 444		x
75  imes 49	$150 \times 222$	X	x
75  imes 49	$225 \times 148$	X	x
75 × 144	$300 \times 324$	Х	x
$100 \times 108$	$100 \times 972$		x
$150 \times 52$	10  imes 7085	Х	
$150 \times 52$	30 × 2361	Х	
$150 \times 52$	75  imes 944	Х	X
$150 \times 72$	$150 \times 648$		X
$150 \times 72$	300 × 324	Х	X
$200 \times 54$	200  imes 486		X
$225 \times 48$	$300 \times 324$	Х	x
$250 \times 43$	250 × 388		X
$100 \times 108$	300 × 324		x
$200 \times 54$	300 × 324		x
$250 \times 52$	300 × 393		x

Table 2: **BFM Simulations of Binary Blends, and Clamped Matrix Chains**: Systems used to study probe and tracer diffusion with monomer number fraction of the probe  $f_p = n_p N_p / (n_p N_p + n_m N_m) = 0.10$ . The first two columns show the length and number of the probe and matrix chains, respectively. An "x" in the third or last columns indicates a simulation of binary blends with CR or with clamped matrix chains, respectively.

MD binary blends		SS binary blends		SS no-CR
$N_p \times n_p$	$N_m \times n_m$	$N_p \times n_p$	$N_m \times n_m$	$N_p \times n_p$
50 × 300	$1000 \times 85$	5  imes 60	$100 \times 17$	5  imes 60
$100 \times 150$	$1000 \times 85$	$10 \times 30$	$100 \times 17$	$10 \times 30$
$200 \times 75$	$1000 \times 85$	$20 \times 15$	$100 \times 17$	20  imes 15
		$35 \times 12$	$100 \times 34$	35  imes 12
		$50 \times 12$	$100 \times 34$	50  imes 12
				$100 \times 17$

Table 3: **MD and Slip Spring Simulations**: Systems  $(N \times n)$ , where N is chain length and n is the number of chains. The subscripts "p" and "m" indicate "probe" and "matrix chains", respectively. The last column corresponds to SS simulations in which CR is switched off.

#### 2 Tabulated Self and Tracer Diffusion Data

Tables 4-7 summarize the diffusion coefficients obtained from the BFM, MD, and SS models. A brief description of simulation and parameter settings are provided in the captions to the tables. The number of beads and the diffusion coefficients obtained from the SS simulations are rescaled and expressed in terms of MD units.

N	n	$D_s (\times 10^{-6})$	Ref. Num.
30	3240	$245\pm3$	4
32	125	257	3
40	100	175	3
50	80	115	3
75	1296	$51.6\pm1$	4
80	50	41.8	3
100	1080	$29.4\pm0.4$	this work
150	648	$11.2\pm0.4$	4
160	25	9.62	3
200	540	$5.9\pm0.2$	this work
225	480	$4.4\pm0.1$	this work
250	432	$3.3\pm0.1$	this work
300	45	2.71	3
300	360	$2.15\pm0.23$	2
315	908	$1.84\pm0.07$	1

Table 4: **BFM Monodisperse**: Systems with  $N = N_m = N_p$  simulated here, and in the literature as indicated by the reference number (as listed at the end of SI) in the last column. *n* is the number of chains used in the simulation.

$N_p$	$N_m$	$D(\times 10^{-6})$	$\hat{D}~(\times 10^{-6})$
75	30	$91.9 \pm 1.5$	$\overline{33.8\pm2.8}$
75	75	$51.6\pm1.0$	$28.8\pm1.9$
75	150	$39.1\pm1.3$	$29.5\pm2.2$
75	225	$31.1\pm1.6$	$28.3\pm2.3$
75	300	$36.2\pm2.5$	$29.5\pm0.9$
100	100		$15.1\pm0.7$
150	10	$57.2\pm0.3$	
150	30	$37.9\pm0.5$	
150	75	$23.4\pm0.5$	$6.5\pm0.6$
150	150	$11.2\pm0.4$	$6.2\pm0.4$
150	300	$8.0\pm0.2$	$6.9\pm0.3$
200	200		$3.0\pm0.2$
225	225		$2.3\pm0.2$
225	300	$3.3\pm0.1$	$2.4\pm0.1$
250	250		$1.9\pm0.2$
300	10	$25.1\pm0.5$	
300	30	$16.0\pm0.2$	
300	75	$8.3\pm0.2$	
300	150	$4.4\pm0.3$	
300	300	$2.2\pm0.2$	$1.4\pm0.1$
100	300		$15.6\pm0.7$
200	300		$3.7\pm0.3$
250	300		$1.9\pm0.2$

Table 5: **BFM Blends**: Systems used to study probe and tracer diffusion with  $\phi_p = 0.05$ . The last column shows results from clamped matrix simulations, which is used to estimate  $D_{\infty}$ . The  $N_p = 300$  probe diffusion diffusivities were previously reported in ref. 4.

$N_p$	$D_s (\times 10^{-4})$	$D_p \times 10^{-4}$
5	$149.9\pm7.4$	
10	$72.6\pm3.6$	
25	$24.9\pm1.3$	
50	$10.2\pm0.5$	$7.1\pm0.4$
100	$3.6\pm0.2$	$2.0\pm0.2$
200	$0.9\pm0.1$	$0.6\pm0.1$

Table 6: **MD Simulations**: Self and probe diffusion coefficients of probe chains,  $D_s$  and  $D_p(N_p, N_m = 1000)$ , obtained from MD simulations.  $D_s$  is calculated by averaging over all chains in the monodisperse melt, while  $D_p$  is obtained from the 15% probe chains blended in long chain matrices of chain length  $N_m = 1000$ .

$N_p$	$D_s (\times 10^{-6})$	$D_p \times 10^{-6}$	$\tilde{D}\times 10^{-6}$
50	$9480.0\pm300.0$	$7650.0\pm300.0$	$7590.0\pm300.0$
100	$3490.0\pm100.0$	$2410.0\pm100.0$	$2310.0\pm100.0$
200	$810.4\pm21.0$	$524.0\pm12.0$	$459.0\pm13.0$
350	$187.6\pm7.0$	$133.0\pm5.0$	$113.0\pm5.0$
500	$71.1\pm3.1$	$56.9\pm2.8$	$48.8\pm3.0$
1000	$12.1\pm0.6$	$12.1\pm1.1$	$9.7\pm1.1$
1500	$4.5\pm0.4$		$4.2\pm0.8$

Table 7: SS Simulations: Diffusion coefficients of probe chains in monodisperse melts,  $D_s$ , binary blends consisting of 15% probe chains in a matrix of long chains,  $D_p$ , and in permanent networks in which CR is switched off,  $\tilde{D}$ , obtained from single-chain slip-spring model simulations. The numbers of beads, length and time scales have all been mapped to those used in the MD simulations. The matrix chain length in the binary blends is mapped to  $N_m = 1000$  of the flexible KG bead-spring chains.

#### **3** Comparison with MCSS and PCN Simulations

In figure 1, we overlay the MCSS (multichain slip spring) and PCN (primitive chain network) calculations of Masubuchi and Uneyama on Figure 4 from the paper.<sup>5</sup> As seen from the plot, the range of the SS simulations performed in this work is similar to the MCSS calculations (dashed red lines), and the agreement between these two sets of calculations is excellent within reported error bars. The PCN series lies slightly above the SS and MCSS simulations, and perhaps shows some early signs of transitioning to a pure reptation  $(D_s \sim Z^{-2})$  regime.



Figure 1: Figure 4 from the manuscript overlaid with data from Masubuchi's MCSS (dashed red) and PCN (dashed blue) simulations.

#### 4 Tracer Diffusivity from Clamped Matrix Simulations

In the BFM simulations, we estimate tracer diffusivity using  $D_{\infty}(N_p) \approx \langle \hat{D}(N_p) \rangle$ . The validity of this assumption hinges on the independence of  $\hat{D}$  and matrix molar mass. We tested this insensitivity by performing a suite of CM simulations (results are shown in tables 4 and 5) at different values of  $N_p$  and  $N_m$ . For example, at  $N_p = 75$ , we performed five different simulations at  $N_m = 30, 75, 150, 225, \text{ and } 300$ . The average across these five simulations,  $\langle \hat{D} \rangle = 2.90 \pm 0.20 \times 10^{-5}$ , was used to normalize  $\hat{D}(N_p = 75, N_m)$ . Similar calculations were performed for all the different probes. Figure 2 plots the normalized CM probe diffusivity for all the samples studied with CM simulations. Normalization by  $\langle \hat{D} \rangle$  enables us to plot data at different  $N_p$  on the same plot. The data are clustered around the expected value of unity over the entire range of molecular weights studied. It demonstrates that  $\hat{D}(N_p, N_m)$  is insensitive to  $N_m$  over the range of molecular weights studied here. This generality gives us the confidence to extrapolate  $\hat{D}$  to the  $N_m \gg N_p$  regime, and hence to estimate  $D_{\infty}$ .



Figure 2: Normalized probe diffusivity from CM simulations for  $N_p$  between 75 and 300, and range of  $N_m$  as shown in table 5. The value of the probe diffusivity  $\hat{D}(N_p, N_m)$  at a particular  $N_p$  is normalized by the averaging over different  $N_m$  simulated  $\langle \hat{D}(N_p) \rangle$ . It shows that probe diffusivity from CM simulations is insensitive to  $N_m$ .

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